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# Reply to "Comment on H.C. Rosu, O. Cornejo-Perez and P. Chen, "Nonsingular parametric oscillators Darboux related to the classical harmonic oscillator"" by A. Schulze-Halberg 

h.C Rosu, ${ }^{1}$ P. Cornejo-PÉrez ${ }^{2}$ and P. Chen ${ }^{3}$<br>${ }^{1}$ IPICyT, Instituto Potosino de Investigacion Cientifica y Tecnologica, Apdo Postal 3-74 Tangamanga, 78231 San Luis Potosí, Mexico<br>${ }^{2}$ Facultad de Ingeniería, Universidad Autónoma de Querétaro, Centro Universitario Cerro de las Campanas, 76010 Santiago de Querétaro, Mexico<br>${ }^{3}$ Leung Center for Cosmology and Particle Astrophysics and Department of Physics and Graduate Institute of Astrophysics<br>National Taiwan University, Taipei, 10617 Taiwan

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Abstract -*** Missing author ***

The comment of Schulze-Halberg [1] on our work [2] stems from the formalism of Darboux covariance of second-order differential operators [3]. We thus recall that two second-order linear differential operators $D_{1}$ and $D_{2}$ are Darboux covariant, or isospectral, if: (i) their factorizations are of the type $D_{1}=L_{1} L_{2}$ and $D_{2}=L_{2} L_{1}$, respectively, and (ii) their eigenfunctions are related through $\psi_{2}=L \psi_{1}$, where $L$ is known as the intertwining operator. The $L_{1,2}$ operators are first order and their non-operatorial part is the logarithmic derivative of an eigenfunction $\psi_{1}$, or a linear combination of eigenfunctions of $D_{1}$. In the simplest constructions of this type, the nodeless ground state of $D_{1}$ is employed. The non-operatorial part of the intertwiner $L$ is also a logarithmic derivative of a function that should be linearly independent with respect to $\psi_{1}$. The Darboux covariance led in the 1980s to the development of the rich area of supersymmetric quantum mechanics [4].

On the other hand, our application is in the area of classical mechanics where the Darboux covariance should be applied at fixed zero 'spectral parameter', where, as shown by Schulze-Halberg, there is a direct relationship between the function used in the intertwiner and $\psi_{1}$. Then, at the formal level, Schulze-Halberg obtains closed form expressions of the solutions of the two Darboux partner equations in terms of the factorization coefficients $\alpha$ and $\beta$. This is an important formal result. However, although it
might look like a simplification, we want to point out that one still should substitute into formulas (10) and (11) our expressions of $\alpha$ and $\beta$ in order to show that the solutions are the ones we obtained. In other words, for equations of given structure, $\alpha$ and $\beta$ can be obtained only through the method explained in our papers $[2,5,6]$ in which the factorizations have been used for convenience since it is not easy to disentangle their expressions directly from the equation. Thus, the results of Schulze-Halberg are of relevance once the two factorization coefficients have been determined and for designed differential equations by giving the expressions of $\alpha$ and $\beta$.

Suppose, for example, that we give $\alpha=x$ and $\beta=1$, then the $\alpha \beta$ designed equation is

$$
\begin{equation*}
\psi^{\prime \prime}+\left(x+\frac{2}{x}\right) \psi^{\prime}+\psi=0 \tag{1}
\end{equation*}
$$

whose general solution calculated with the formula (10) of Schulze-Halberg is

$$
\begin{equation*}
\psi=\frac{\sqrt{2 \pi}}{2 x}\left[1+\operatorname{erf}\left(\frac{x}{\sqrt{2}}\right)\right] . \tag{2}
\end{equation*}
$$

On the other hand, the Darboux partner equation is

$$
\begin{equation*}
\phi^{\prime \prime}+x \phi^{\prime}+\phi=0 \tag{3}
\end{equation*}
$$

whose general solution calculated with formula (11) of the Comment is

$$
\begin{equation*}
\phi=\exp \left(-\frac{x^{2}}{2}\right) \int^{x} \exp \left(\frac{t^{2}}{2}\right) d t \tag{4}
\end{equation*}
$$

However, if the equation to be analysed is

$$
\begin{equation*}
y^{\prime \prime}+\frac{4 \omega_{0}}{\sin 2 \omega_{0} t} y^{\prime}+\omega_{0}^{2} y=0 \tag{5}
\end{equation*}
$$

and one knows a priori it is an $\alpha \beta$-designed equation but without knowing $\alpha$ and $\beta$, it will not be easy to find $\alpha$ and $\beta$ and one should use standard methods such as those in our papers. In fact, equation (5) corresponds to the case $\lambda=1$ and therefore it is still a singular oscillator equation but we use it herein for illustrative purposes. Then, it is easy to check that

$$
\begin{equation*}
y=\frac{\omega_{0}}{\sin \omega_{0} t} \tag{6}
\end{equation*}
$$

is a particular solution of (5) that we extracted from the expression of the general solution given in our paper [2] for $C_{1}=0$ and $C_{2}=\omega_{0}$.

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