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# Itinerary synchronization between PWL systems coupled with unidirectional links

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## Abstract

In this paper the collective dynamics of  $N$ -coupled piecewise linear (PWL) systems with different number of scrolls is studied. The coupling is in a master-slave sequence configuration, with this type of coupling we investigate the synchrony behavior of a ring-connected network and a chain-connected network both with unidirectional links. Itinerary synchronization is used to detect synchrony behavior. Itinerary synchronization is defined in terms of the symbolic dynamics arising by assigning different numbers to the regions where the scrolls are generated. A weaker variant of this notion,  $\epsilon$ -itinerary synchronization is also introduced and numerically investigated. It is shown that in certain parameter regimes if the inner connection between nodes takes account of all the state variables of the system (by which we mean that the inner coupling matrix is the identity matrix), then itinerary synchronization occurs and the coordinate motion is determined by the node with the smallest number of scrolls. Thus the collective behavior in all the nodes of the network is determined by the node with least scrolls in its attractor. Results about the dynamics in a directed chain topology are also presented. Depending on the inner connection properties, the nodes present multistability or preservation of the number of scrolls of the attractors.

**keywords:** Itinerary Synchronization; chaos; dynamical networks; multiscroll attractor.

## 23 1 Introduction

24 Piecewise linear (PWL) systems are used to construct simple chaotic oscillators  
25 capable of generating various multiscroll attractors in the phase space. These  
26 systems contain a linear part plus a nonlinear element characterized by a switch-  
27 ing law. One of the most studied PWL system is the so called Chua's circuit,  
28 whose nonlinear part (also named the Chua diode) generates two scroll attrac-  
29 tors [1, 2]. Inspired by the Chua circuit, a great number of PWL systems have  
30 been produced via various switching systems [3]. A review and summary of  
31 different approaches to generate multiscroll attractors can be found in [4, 5, 6]  
32 and references therein.

33 Synchronization phenomena in a pair of coupled PWL systems has also at-  
34 tracted attention in the context of nonlinear dynamical systems theory and its  
35 applications [7, 8].

36 In general, we say that a set of dynamical systems achieve synchronization  
37 if trajectories in each system approach a common trajectory (in some sense) by  
38 means of interactions [9].

39 One way to study synchronization in a pair of PWL systems is to couple  
40 them in a master-slave configuration [1, 10]. In [11] the dynamical mechanism  
41 leading to projective synchronization of Chua circuits with different scrolls is  
42 investigated. In [12], a master-slave system composed of PWL systems is con-  
43 sidered in which the slave system displays more scrolls in its attractor than  
44 the master system. The main result is that the slave system synchronizes with  
45 the master system by reducing its number of attractor scrolls, while the master  
46 preserves its number of scrolls. A consequence is the emergence of multistabil-  
47 ity phenomena. For instance, if the number of scrolls presented by the master  
48 system is less than the number of scrolls presented by the slave system, then  
49 the slave system can oscillate in multiple basins of attraction depending on its  
50 initial condition. Conversely, when the system of [12] is adjusted so that the  
51 master system displays more scrolls than the slave system when uncoupled then  
52 the slave system increases its number of attractor scrolls to equal that of the  
53 master system when coupled.

54 We study a system composed of an ensemble of master-slave systems coupled  
55 in a ring configuration network; *i.e.*, a dynamical network where each node is  
56 a PWL-system with varying numbers of scrolls in the attractors and connected  
57 in a ring topology with directional links. In order to address this problem, we  
58 introduce three concepts: 1) scroll-degree, which is defined as the number of  
59 scrolls of an attractor in a given node; 2) a network of nearly identical nodes,  
60 *i.e.*, a dynamical network composed of PWL systems with perhaps different  
61 scroll degree but similar underlying differential equations and 3) itinerary syn-  
62 chronization based on symbolic dynamics. A PWL system is defined by means  
63 of a partition of the space where linear systems act, so this natural partition is  
64 useful for analyzing synchronization between dynamical systems by using sym-  
65 bolic dynamics. Of course itinerary synchronization does not imply complete  
66 synchronization, where trajectories converge to a single one. In this paper we  
67 study the emergence of itinerary synchronization,  $\epsilon$ -itinerary synchronization,

68 multistability and the preservation of the scroll number of a network of nearly  
69 identical nodes. Furthermore, we remove a link to the ring topology in order  
70 to study the effect of topology changes in the collective dynamics of the net-  
71 work. That is, we modify the topology by deleting a single link, transforming  
72 the structure to a directed chain of coupled systems which we call an open ring.  
73 We have formulated two possible scenarios after the link deletion: a) the first  
74 node in the chain has the largest scroll-degree or, b) it has the smallest one. In  
75 both scenarios we assume that the inner coupling matrix is the identity matrix  
76 *i.e.* the coupling between any pair of nodes is throughout all its state variables.

77 To the best of our knowledge, multistability and scroll-degree preservation  
78 have not been studied in the context of PWL dynamical networks. We note that  
79 Zhao *et.al.* in [14] established synchronization criteria for certain networks of  
80 non-identical nodes with the same equilibria point [14]. The authors proposed  
81 stability conditions in terms of inequalities involving matrix spectra which are,  
82 computationally speaking, difficult to solve. Sun *et.al.* in [13] studied the  
83 case in which nodes are nearly identical in the sense that each node has a slight  
84 parametric mismatch. The authors proposed an extension of the master stability  
85 functions for these types of dynamical network.

86 We have organized this paper as follows: In section 2 we introduce some  
87 mathematical preliminaries. In section 3 we give an easy approach to generate  
88 a one dimensional grid multiscroll attractor via PWL systems. In section 4 we  
89 introduce a partition to configure the symbolic dynamics of trajectories of a pair  
90 of coupled PWL systems. In section 5 we propose a definition of itinerary syn-  
91 chronization based on the itinerary of trajectories of a master-slave system. In  
92 section 6 we give some preliminaries of dynamical networks which are composed  
93 of  $N$  coupled dynamical systems. In section 7 the dynamics of  $N$ -coupled PWL  
94 systems in a ring topology network is analyzed. Some examples about itinerary  
95 synchronization are studied and different forms of couplings are also considered.  
96 Finally, in section 8 we discuss conclusions.

## 97 2 Mathematical Preliminaries

### 98 2.1 Piecewise linear dynamical systems

99 Let  $T : X \rightarrow X$ , with  $X \subset \mathbb{R}^n$  and  $n \in \mathbb{Z}^+$ , be a piecewise linear dynamical  
100 system whose dynamics is given by a family of sub-systems of the form

$$\dot{\mathcal{X}} = A_\tau \mathcal{X} + B_\tau, \quad (1)$$

101 where  $\mathcal{X} = (x_1, \dots, x_n)^T \in \mathbb{R}^n$  is the state vector,  $A_\tau = \{\alpha_{ij}^\tau\} \in \mathbb{R}^{n \times n}$ , with  
102  $\alpha_{ij}^\tau \in \mathbb{R}^+$ , and  $B_\tau = (\beta_{\tau 1}, \dots, \beta_{\tau n})^T \in \mathbb{R}^n$  are the linear operators and constant  
103 real vectors of the  $\tau$ th-subsystems, respectively. The index  $\tau \in \mathcal{I} = \{1, \dots, \eta\}$  is  
104 given by a rule that switches the activation of a sub-system in order to determine  
105 the dynamics of the PWL system. Let  $X$  be a subset of  $\mathbb{R}^n$  and  $\mathcal{P} = \{P_1, \dots, P_\eta\}$   
106 ( $\eta > 1$ ) be a finite partition of  $X$ , that is,  $X = \bigcup_{1 \leq i \leq \eta} P_i$ , and  $P_i \cap P_j = \emptyset$  for  
107  $i \neq j$ . Each element of the set  $\mathcal{P}$  is called an atom.

108 The selection of the index  $\tau$  can be given according to a predefined itinerary  
 109 and controlling by time; or by requiring that  $\tau$  takes its value according to  
 110 the state variable  $\chi$  depending upon which atom of a finite partition of the  
 111 state-space  $\mathcal{P} = \{P_1, \dots, P_\eta\}$  ( $\eta \in \mathbb{Z}^+$ ) a point is in.

112 An easy way to generate a partition  $\mathcal{P}$  is to consider a vector  $\mathbf{v} \in \mathbb{R}^n$  (with  
 113  $\mathbf{v} \neq 0$ ) and a set of scalars  $\delta_1 < \delta_2 < \dots < \delta_{\eta-1}$  such that each  $P_i = \{\mathcal{X} \in$   
 114  $\mathbb{R}^n : \delta_{i-1} \leq \mathbf{v}^T \mathcal{X} < \delta_i\}$ , with  $i = 2, \dots, \eta - 1$ ,  $P_1 = \{\mathcal{X} \in \mathbb{R}^n : \mathbf{v}^T \mathcal{X} < \delta_1\}$ ,  
 115 and  $P_\eta = \{\mathcal{X} \in \mathbb{R}^n : \delta_{\eta-1} \leq \mathbf{v}^T \mathcal{X}\}$ . We call the hyperplanes  $\mathbf{v}^T \mathcal{X} = \delta_i$   
 116 ( $i = 1, \dots, \eta - 1$ ) the switching surfaces. Without loss of generality, we assume  
 117 that the hyperplanes  $\mathbf{v}^T \mathcal{X} = \delta_i$  (for  $i = 1, 2, \dots, \eta - 1$ ) are defined with  $\mathbf{v} =$   
 118  $(1, 0, \dots, 0)^T \in \mathbb{R}^n$ .

119 In this paper we consider a piecewise linear system  $(T, \mathcal{P})$ , such that its  
 120 restriction to each atom  $P_i$  has a fixed point  $\mathcal{X}_i^*$  i.e.  $T(\mathcal{X}_i^*) = 0$  for one  $\mathcal{X}_i^* \in P_i$   
 121 ( $i \in \mathcal{I}$ ). Clearly  $\mathcal{X}_i^* = -A_\tau^{-1} B_\tau$ . We assume that the switching signal depends  
 122 on the state variable and is defined as follows:

123 *Definition 2.1.* Let  $\mathcal{I} = \{1, 2, \dots, \eta\}$  be an index set that labels each element  
 124 of the family of the sub-systems (1). A function  $\kappa : \mathbb{R}^n \rightarrow \mathcal{I} = \{1, \dots, \eta\}$  of the  
 125 form

$$\kappa(\mathcal{X}) = \begin{cases} 1, & \text{if } \mathcal{X} \in P_1; \\ 2, & \text{if } \mathcal{X} \in P_2; \\ \vdots & \vdots \\ \eta, & \text{if } \mathcal{X} \in P_\eta; \end{cases} \quad (2)$$

126 is called a switching signal. Furthermore, if  $\kappa(\mathcal{X}) = \tau_i \in \mathcal{I}$  is the value  
 127 of the switching signal during the time interval  $t \in [t_i, t_{i+1})$ , then  $\mathcal{S}(\mathcal{X}_0) =$   
 128  $\{\tau_0, \tau_1, \dots, \tau_m, \dots\}$  gives the itinerary generated by  $\kappa(\mathcal{X}_0)$  at  $\mathcal{X}_0$  and,  $\mathcal{S}(i, \mathcal{X}_0)$   
 129 is the element  $\tau_i \in \mathcal{S}(\mathcal{X}_0)$  that occurs at time  $t_i$ , this defines a set of switching  
 130 times  $\Delta_t = \{t_0, t_1, \dots, t_m, \dots\}$ .

131 Note that  $\tau$  changes only when the orbit  $\phi(t, \chi_0)$  goes from one atom  $P_i$  to  
 132 another  $P_j$ ,  $i \neq j$ .

133 *Definition 2.2.* A  $\eta$ -PWL system is composed of two sets:  $\mathbf{A} = \{A_1, \dots, A_\eta\}$   
 134 and  $\mathbf{B} = \{B_1, B_2, \dots, B_\eta\}$ , with  $A_\tau = \{\alpha_{ij}^\tau\} \in \mathbb{R}^{n \times n}$  ( $\alpha_{ij}^\tau \in \mathbb{R}$ ) and  $B_\tau =$   
 135  $(\beta_{\tau 1}, \dots, \beta_{\tau n})^T \in \mathbb{R}^n$ ; and a switching signal  $\kappa : \mathbb{R}^n \rightarrow \mathcal{I} = \{1, 2, \dots, \eta\}$  so that:

$$\dot{\mathcal{X}} = \begin{cases} A_1 \mathcal{X} + B_1, & \text{if } \kappa(\mathcal{X}) = 1; \\ A_2 \mathcal{X} + B_2, & \text{if } \kappa(\mathcal{X}) = 2; \\ \vdots & \vdots \\ A_\eta \mathcal{X} + B_\eta, & \text{if } \kappa(\mathcal{X}) = \eta. \end{cases} \quad (3)$$

136 We can rewrite (3) in a more compact form as:

$$\dot{\mathcal{X}} = A_{\kappa(\mathcal{X})} \mathcal{X} + B_{\kappa(\mathcal{X})}. \quad (4)$$

137 *Definition 2.3.* Two  $\eta_1$ -PWL and  $\eta_2$ -PWL systems are called quasi-symmetrical  
 138 if they are governed by the same linear operator  $A = A_i$  for all  $i$  but  $\eta_1 \neq \eta_2$ .

139 **3 System Description: one direction grid scrolls**  
 140 **attractor**

141 Now we assume that the dimension of each  $\eta$ -PWL system is  $n = 3$  and that  
 142 the eigenspectra of linear operators  $A_\tau \in \mathbb{R}^{3 \times 3}$  have the following features: a)  
 143 one eigenvalue is a real number; and b) two eigenvalues are complex conjugate  
 144 numbers with non-zero imaginary part. There is an approach to generate dy-  
 145 namical systems based on these linear dissipative systems in the case where the  
 146 complex eigenvalues and the real eigenvalue have mixed sign (sometimes called  
 147 an unstable dissipative system (UDS) [15]). In this paper we use a particular  
 148 type of unstable dissipative system (UDS) called *Type I*:

149 *Definition 3.1.* A subsystem  $(A_\tau, B_\tau)$  of the system (4) in  $\mathbb{R}^3$  is said to be an  
 150 UDS of *Type I* if the eigenvalues of the linear operator  $A_\tau$  denoted by  $\lambda_i$  satisfy:  
 151  $\sum_{i=1}^3 \lambda_i < 0$ ;  $\lambda_1$  is a negative real eigenvalue and; the other two  $\lambda_2$  and  $\lambda_3$  are  
 152 complex conjugate eigenvalues with positive real part. The system is an UDS  
 153 of *Type II* if  $\sum_{i=1}^3 \lambda_i < 0$ , and one  $\lambda_i$  is a positive real eigenvalue and; the other  
 154 two  $\lambda_i$  are complex conjugate eigenvalues with negative real part.

155 To each  $\tau \in \mathcal{I}$  is associated an atom  $P_\tau \subset \mathbb{R}^n$ , containing an equilib-  
 156 rium point  $\chi_\tau^* = -A^{-1}B_\tau$  which has a one-dimensional stable manifold  $E^s =$   
 157  $Span\{\bar{v}_j \in \mathbb{R}^3 : \alpha_j < 0\}$  and a two-dimensional unstable manifold  $E^u =$   
 158  $Span\{\bar{v}_j \in \mathbb{R}^3 : \alpha_j > 0\}$ , with  $\bar{v}_j$  an eigenvector of the linear operator  $A$  and  
 159  $\lambda_j = \alpha_j + i\beta_j$  its corresponding eigenvalue; *i.e.* it is a saddle equilibrium point.  
 160 We are interested in bounded flows which are generated by quasi-symmetrical  
 161  $\eta$ -PWL systems such that for any initial condition  $\mathcal{X}_0 \in \mathbb{R}^3$ , the orbit  $\phi(t, \chi_0)$   
 162 of the  $\eta$ -PWL system (4) limits to a one-spiral trajectory in the atom  $P_\tau$  called  
 163 a scroll. The orbit escapes from one atom to other due to the unstable mani-  
 164 fold in each atom. In this context, the system  $\eta$ -PWL (4) can display various  
 165 multi-scroll attractors as a result of a combination of several unstable one-spiral  
 166 trajectories, while the switching between regions is governed by the function  
 167 (2).

168 *Definition 3.2.* The scroll-degree of a  $\eta$ -PWL system (4) based on UDS *Type*  
 169 *I* is the maximum number of scrolls that the PWL system can display in the  
 170 attractor.

171 In this work we consider the same linear operator  $A$ , so  $A_\tau = A$  for all  $\tau$ .  
 172 An easy approach to generate a one dimensional grid multiscroll attractor via a  
 173 PWL system based on UDS type I form is by defining a double-scroll attractor  
 174 as follows:

- 175 • Consider the linear operator  $A$ :

$$A = \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ -\alpha_{31} & -\alpha_{32} & -\alpha_{33} \end{pmatrix}, \quad (5)$$

176 where  $\alpha_{31}$ ,  $\alpha_{32}$  and  $\alpha_{33}$  satisfy the UDS type I conditions, *i.e.*,  $\lambda_1 \in \mathbb{R}$ , and

177  $\lambda_2, \lambda_3 \in \mathbb{C}$  such that the absolute value of the imaginary part is greater  
 178 than the absolute value of the real part of  $\lambda_i$ , with  $i = 2, 3$ .

- 179 • Choose two equilibria on the  $x$ -axis:  $\chi_1^* = (x_{eq1}^*, 0, 0)^T$  and  
 180  $\chi_2^* = (x_{eq2}^*, 0, 0)^T$ .
- 181 • Compute the stable and unstable manifolds  $E_1^s, E_1^u, E_2^s$ , and  $E_2^u$  associated  
 182 to each equilibria  $\chi_1^*$  and  $\chi_2^*$ , respectively.
- 183 • Find the intersection points between the stable manifold  $E_1^s$  and the un-  
 184 stable manifold  $E_2^u$ , and between the stable manifold  $E_2^s$  and the unstable  
 185 manifold  $E_1^u$ .
- 186 • Define the switching surface as the plane that pass through the intersection  
 187 points  $E_1^s \cap E_2^u$ , and  $E_2^s \cap E_1^u$  and the line:  $x_1 = (x_{eq1}^* + x_{eq2}^*)/2, x_3 = 0$ .
- 188 • Compute the constant vectors  $B_\tau = -A\chi_\tau^*$ , with  $\tau = 1, 2$ .

189 The above steps generate two heteroclinic orbits between the equilibria  $\chi_1^*$   
 190 and  $\chi_2^*$ . One of the heteroclinic orbits is from  $\chi_1^*$  to the point  $E_2^s \cap E_1^u$  and from  
 191 this point to  $\chi_2^*$ . The other heteroclinic orbit is from  $\chi_2^*$  to the point  $E_1^s \cap E_2^u$   
 192 and from this point to  $\chi_1^*$ .

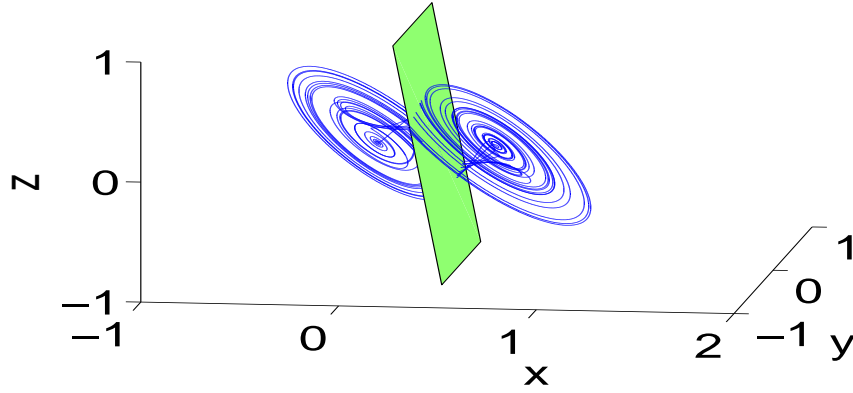


Figure 1: Attractor generated by the PWL system given by the linear operator (5), The vectors  $B_1 = (0, 0, 0)^T$  and  $B_2 = (0, 0, 0.9)^T$  and switching surface  $\{\mathcal{X} \in \mathbb{R}^3 : 0.7369x_1 + 0.0918x_3 - 0.2211 = 0\}$  (green plane).

193 In order to illustrate the approach to generate double-scroll attractors using  
 194 (4), we set  $\alpha_{31} = 1.5, \alpha_{32} = 1$  and  $\alpha_{33} = 1$ .

- 195 • Thus, the eigenvalues are  $\lambda_1 = -1882/1563, \lambda_2 = 319/3126 + 2503/2252i,$   
 196 and  $\lambda_3 = 319/3126 - 2503/2252i$  which satisfy:  $\sum_{i=1}^3 \lambda_i < 0$  and  
 197  $Imag(\lambda_2)/Re(\lambda_2) > 6$ .  $Imag(\lambda_2)$  and  $Re(\lambda_2)$  denote the imaginary part  
 198 and real part of  $\lambda_2$ , respectively.

- 199 • Choose equilibria at  $\chi_1^* = (0, 0, 0)^T$  and  $\chi_2^* = (0.6, 0, 0)^T$ .
- 200 • The unstable manifolds  $E_1^u = \{\mathcal{X} \in \mathbb{R}^3 : 0.3646x_1 - 0.0597x_2 + 0.2927x_3 =$   
201  $0\}$  and  $E_2^u = \{\mathcal{X} \in \mathbb{R}^3 : 0.3646x_1 - 0.0597x_2 + 0.2927x_3 - 0.2188 = 0\}$   
202 and the stable manifolds  $E_1^s = \{\mathcal{X} \in \mathbb{R}^3 : \frac{x_1}{-0.4687} = \frac{x_2}{0.5644} = \frac{x_3}{-0.6796}\}$  and  
203  $E_2^s = \{\mathcal{X} \in \mathbb{R}^3 : \frac{x_1 - 0.6}{-0.4687} = \frac{x_2}{0.5644} = \frac{x_3}{-0.6796}\}$
- 204 •  $E_1^s \cap E_2^u = (0.2541, -0.3060, 0.3684)^T$ , and  
205  $E_2^s \cap E_1^u = (0.3459, -0.3060, -0.3684)^T$  and the line:  $x_1 = 0.3, x_2 \in \mathbb{R},$   
206  $x_3 = 0$ . So the switching surface is given by  $\{\mathcal{X} \in \mathbb{R}^3 : 0.7369x_1 +$   
207  $0.0918x_3 - 0.2211 = 0\}$ .
- 208 •  $B_1 = -A\chi_1^* = (0, 0, 0)^T$  and  $B_2 = -A\chi_2^* = (0, 0, 0.9)^T$ .

209 The calculated values approximate the exact values needed for the hetero-  
210 clinic orbit and they allow us to generate a double-scroll attractor by trapping  
211 the trajectories oscillating around the equilibria, see Figure 1.

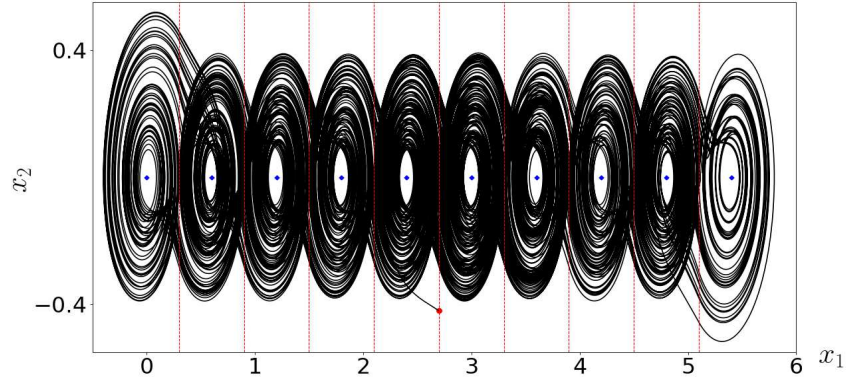


Figure 2: Projection of the attractor generated by the quasi-symmetrical 10-PWL(S) system onto the  $(x_1, x_2)$  plane. The dashed lines mark the division between the atoms.

212 **Example 3.3.** In order to illustrate the generation of multiscroll attractors  
213 using (4), we consider a quasi-symmetrical 10-PWL system defined in  $\mathbb{R}^3$  with  
214 state vector  $\mathcal{X} = (x_1, x_2, x_3)^T$  and linear operator defined as follows

$$A = \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ -\alpha_{31} & -\alpha_{32} & -\alpha_{33} \end{pmatrix}; \quad (6)$$

where  $\alpha_{31} = 1.5, \alpha_{32} = 1$  and  $\alpha_{33} = 1$ ; the set of constants vectors

$$\mathbf{B} = \{B_1 = (0, 0, 0)^T, B_2 = (0, 0, 0.9)^T, B_3 = (0, 0, 1.8)^T, B_4 = (0, 0, 2.7)^T\},$$



$$B_5 = (0, 0, 3.6)^T, B_6 = (0, 0, 4.5)^T, B_7 = (0, 0, 5.4)^T, B_8 = (0, 0, 6.3)^T, \\ B_9 = (0, 0, 7.2)^T, B_{10} = (0, 0, 8.1)^T};$$

215 and the partition:

$$\mathcal{P} = \{ P_1 = \{\mathcal{X} \in \mathbf{R}^3 : x_1 < 0.3\}, P_2 = \{\mathcal{X} \in \{\mathcal{X} \in \mathbf{R}^3 : 0.3 \leq x_1 < 0.9\}, \\ P_3 = \{\mathcal{X} \in \mathbf{R}^3 : 0.9 \leq x_1 < 1.5\}, P_4 = \{\mathcal{X} \in \mathbf{R}^3 : 1.5 \leq x_1 < 2.1\}, \\ P_5 = \{\mathcal{X} \in \mathbf{R}^3 : 2.1 \leq x_1 < 2.7\}, P_6 = \{\mathcal{X} \in \mathbf{R}^3 : 2.7 \leq x_1 < 3.3\}, \\ P_7 = \{\mathcal{X} \in \mathbf{R}^3 : 3.3 \leq x_1 < 3.9\}, P_8 = \{\mathcal{X} \in \mathbf{R}^3 : 3.9 \leq x_1 < 4.5\} \\ P_9 = \{\mathcal{X} \in \mathbf{R}^3 : 4.5 \leq x_1 < 5.1\}, P_{10} = \{\mathcal{X} \in \mathbf{R}^3 : x_1 \geq 5.1\} \} \quad (7)$$

216 The eigenvalues of  $A$  are  $\lambda_1 = -1.20$  and  $\lambda_{2,3} = 0.10 \pm 1.11i$ . By Definition  
217 2.4, the system is an UDS of *Type I*. The equilibrium points for this system are  
218 at  $\chi_1^* = (0, 0, 0)^T$ ,  $\chi_2^* = (0.6, 0, 0)^T$ ,  $\chi_3^* = (1.2, 0, 0)^T$ ,  $\chi_4^* = (1.8, 0, 0)^T$ ,  $\chi_5^* =$   
219  $(2.4, 0, 0)^T$ ,  $\chi_6^* = (3, 0, 0)^T$ ,  $\chi_7^* = (3.6, 0, 0)^T$ ,  $\chi_8^* = (4.2, 0, 0)^T$ ,  $\chi_9^* = (4.8, 0, 0)^T$   
220 and  $\chi_{10}^* = (5.4, 0, 0)^T$ . Figure (2) depicts the projection of the attractor gener-  
221 ated by the quasi-symmetrical 10-PWL(S) system onto the  $(x_1, x_2)$  plane with  
222 initial condition  $\chi_0 = (2.7, -0.42, 0.09)^T$ . We solved this system (3) numerically  
223 by using fourth order Runge-Kutta method with 2,000,000 time iterations and  
224 step-size  $h = 0.01$  in order to corroborate that the system always oscillates in  
225 the attractor and for the initial condition considered the asymptotic regime is  
226 achieved after 5000 iterations. When we refer to 2000 arbitrary units of time  
227 correspond to 200,000 iterations.

The trajectory  $\mathcal{X}(t)$  of the PWL system can be calculated by  $\mathcal{X}^i(t) = e^{At} \mathcal{X}_0^i$   
in each atom  $P_i$ , where  $\mathcal{X}^i = \mathcal{X} + \mathcal{X}_i^*$  and  $\mathcal{X}_0^i$  is the initial condition when the  
trajectory enter to the atom  $P_i$ ,  $i = 1, \dots, 10$ . Then

$$\mathcal{X}^i(t) = PE(t)P^{-1}\mathcal{X}^i(0),$$

228 where  $P$  is the invertible matrix defined by the eigenvector of  $A$  and

$$E(t) = \begin{pmatrix} e^{\lambda_1 t} & 0 & 0 \\ 0 & e^{Re(\lambda_2)t} \sin(Imag(\lambda_2)t) & -e^{Re(\lambda_2)t} \cos(Imag(\lambda_2)t) \\ 0 & e^{Re(\lambda_2)t} \cos(Imag(\lambda_2)t) & e^{Re(\lambda_2)t} \sin(Imag(\lambda_2)t) \end{pmatrix}.$$

## 229 4 Symbolic dynamics of trajectories of a pair of 230 coupled PWL systems

231 Consider a pair of quasi-symmetrical  $\eta$ -PWL systems defined by (4), *i.e.*, they  
232 have different scroll-degrees. They are coupled in a Master-Slave configuration  
233 as follows.

$$\dot{\mathcal{X}}_m = A\mathcal{X}_m + B_{\kappa_m}(\mathcal{X}_m), \\ \dot{\mathcal{X}}_s = A\mathcal{X}_s + B_{\kappa_s}(\mathcal{X}_s) + c\Gamma(\mathcal{X}_m - \mathcal{X}_s), \quad (8)$$

234 where  $\mathcal{X}_m = (x_1^m, x_2^m, x_3^m)^T$  and  $\mathcal{X}_s = (x_1^s, x_2^s, x_3^s)^T$  are the state vectors of  
235 the master and slave systems, respectively.  $\mathcal{X}_m$  is in the phase space of the

236 master system  $D_m$ ;  $\mathcal{X}_s$  is in the phase space  $D_s$ . Clearly the orbits of the  
 237 overall system lie in a subspace of the whole state space  $D_m \oplus D_s$ .  $\kappa_i : \mathbb{R}^3 \rightarrow$   
 238  $\mathcal{I}_i = \{1, 2, \dots, \eta_i\}$ , with  $i = m, s$  and  $\eta_m \neq \eta_s$ , is the signal of the master  
 239 system ( $i = m$ ) and slave system ( $i = s$ ). The itineraries generated by  $\tau$  of the  
 240 master and slave systems are  $\mathcal{S}_m(\mathcal{X}_{m0}) = \{\tau_0, \tau_1, \dots\}$  and  $\mathcal{S}_s(\mathcal{X}_{s0}) = \{\tau'_0, \tau'_1, \dots\}$ ,  
 241 respectively. The corresponding time sets are given by  $\Delta_{tm} = \{t_0, t_1, \dots\}$  and  
 242  $\Delta_{ts} = \{t'_0, t'_1, \dots\}$ . We take the constant matrix  $\Gamma = \text{diag}\{r_1, r_2, r_3\} \in \mathbb{R}^{3 \times 3}$   
 243 to be the inner linking matrix where  $r_l = 1$  (for  $l = 1, 2, 3$ ) if both master and  
 244 slave systems are linked through their  $l$ -th state variable, and  $r_l = 0$  otherwise.  
 245 The parameter  $0 < c \in \mathbb{R}$  is the coupling strength.

246 There are several definitions of synchronization [16, 18], for instance, complete  
 247 synchronization is given as follows:

248 *Definition 4.1.* The master-slave system (4) is said to achieve complete synchro-  
 249 nization if

$$\lim_{t \rightarrow \infty} \|\phi_m(t, \mathcal{X}_{m0}) - \phi_s(t, \mathcal{X}_{s0})\| \rightarrow 0. \quad (9)$$

250 for all initial conditions,  $\mathcal{X}_{m0}$  and  $\mathcal{X}_{s0}$ .

251 The symbol  $\|\cdot\|$  denotes the Euclidean distance in  $\mathbb{R}^3$ . This mode of syn-  
 252 chronization is very strong. There are weaker and more generalized notions of  
 253 synchronization [17]. Suppose that  $\mathcal{F}$  is a transformation from the trajectories  
 254 of the attractor in  $D_m$  space to the trajectories in  $D_s$  space. The precise form  
 255 of  $\mathcal{F}$  will depend upon the application in mind. Given such a transformation,  
 256 Generalized Synchronization is defined as follows.

257 *Definition 4.2.* The master-slave system (4) is said to achieve generalized syn-  
 258 chronization if

$$\lim_{t \rightarrow \infty} \|\mathcal{F}(\phi_m(t, \mathcal{X}_{m0})) - \phi_s(t, \mathcal{X}_{s0})\| \rightarrow 0. \quad (10)$$

259 for all  $\mathcal{X}_{m0}$  and  $\mathcal{X}_{s0}$  where  $\mathcal{F}$  is the given transformation from the trajectories  
 260 of the attractor in  $D_m$  space to the trajectories in  $D_s$  space.

261 It has been reported in [12] that in the type of configuration given by (8) the  
 262 master system determines the scroll-degree in the slave system. In particular, if  
 263  $\eta_m < \eta_s$ , then the master-slave system achieves generalized synchronization and  
 264  $\eta_s - \eta_m + 1$  different basins of attraction appear. The trajectories of the slave  
 265 system depend on their initial condition. That is, the master-slave configuration  
 266 results in multiple basins of attraction for the slave. This phenomenon is called  
 267 multistability [19]. On the other hand, if  $\eta_m > \eta_s$ , then the slave system  
 268 increases its scroll-degree till it matches the master's scroll-degree.

269 In order to illustrate the dynamical behavior of the master-slave system,  
 270 consider two quasi-symmetrical  $\eta$ -PWL systems with common linear operator  
 271  $A$  and a set of constant vectors  $\mathbf{B} = \{B_3, B_4, \dots, B_{10}\}$  defined in Example 3.3  
 272 (Eq. (6)).

273 **Example 4.3.** As a first example of a coupled pair of multiscroll chaotic sys-  
 274 tems, suppose that the master's scroll-degree is  $\eta_m = 3$  and the slave's scroll-  
 275 degree is  $\eta_s = 8$ , and both are connected with a coupling strength  $c$  and an

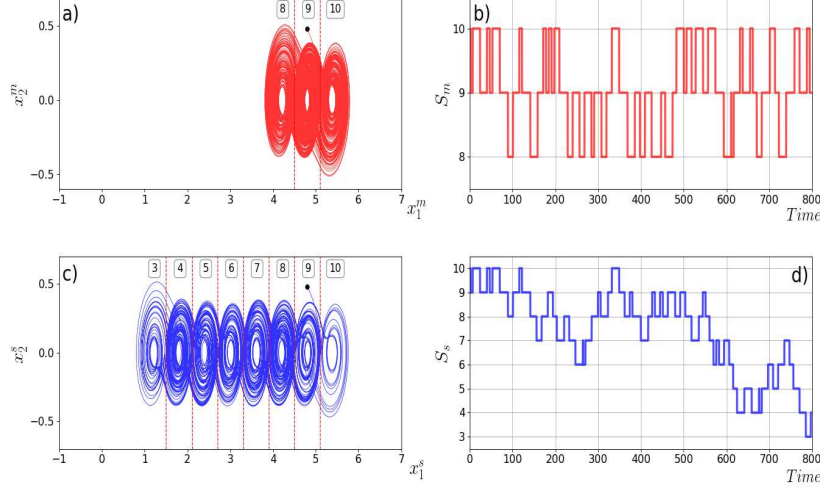


Figure 3: a) Projection of the master system onto the plane  $(x_1^m, x_2^m)$  with initial condition  $\chi_{m0} = (4.8, 0.48, -0.29)^T$ ; b) The master itinerary  $S_m(\chi_{m0})$ ; c) Projection of the slave system onto the plane  $(x_1^s, x_2^s)$  with initial condition  $\chi_{s0} = (4.8, 0.48, -0.29)^T$ , for coupling strength  $c = 0$ ; d) The slave itinerary  $S_s(\chi_{s0})$ .

276 inner coupling matrix given by  $\Gamma = \{0, 1, 0\}$ . The signal for the master system  
 277  $\kappa_m : \mathbb{R}^3 \rightarrow \mathcal{I}_m = \{8, 9, 10\}$  is

$$\kappa_m(\mathcal{X}) = \begin{cases} 10, & \text{if } \mathcal{X} \in P_{10} = \{\mathcal{X} \in \mathbb{R}^3 : x_1 \geq 5.1\}; \\ 9, & \text{if } \mathcal{X} \in P_9 = \{\mathcal{X} \in \mathbb{R}^3 : 4.5 \leq x_1 < 5.1\}; \\ 8, & \text{if } \mathcal{X} \in P_8 = \{\mathcal{X} \in \mathbb{R}^3 : x_1 < 4.5\}. \end{cases} \quad (11)$$

278 And for the slave system the function  $\kappa_s : \mathbb{R}^3 \rightarrow \mathcal{I}_s = \{3, 4, \dots, 10\}$  is

$$\kappa_s(\mathcal{X}) = \begin{cases} 10, & \text{if } \mathcal{X} \in P_{10} = \{\mathcal{X} \in \mathbb{R}^3 : x_1 \geq 5.1\}; \\ 9, & \text{if } \mathcal{X} \in P_9 = \{\mathcal{X} \in \mathbb{R}^3 : 4.5 \leq x_1 < 5.1\}; \\ 8, & \text{if } \mathcal{X} \in P_8 = \{\mathcal{X} \in \mathbb{R}^3 : 3.9 \leq x_1 < 4.5\}; \\ 7, & \text{if } \mathcal{X} \in P_7 = \{\mathcal{X} \in \mathbb{R}^3 : 3.3 \leq x_1 < 3.9\}; \\ 6, & \text{if } \mathcal{X} \in P_6 = \{\mathcal{X} \in \mathbb{R}^3 : 2.7 \leq x_1 < 3.3\}; \\ 5, & \text{if } \mathcal{X} \in P_5 = \{\mathcal{X} \in \mathbb{R}^3 : 2.1 \leq x_1 < 2.7\}; \\ 4, & \text{if } \mathcal{X} \in P_4 = \{\mathcal{X} \in \mathbb{R}^3 : 1.5 \leq x_1 < 2.1\}; \\ 3, & \text{if } \mathcal{X} \in P_3 = \{\mathcal{X} \in \mathbb{R}^3 : x_1 < 1.5\}. \end{cases} \quad (12)$$

279 Using Runge-Kutta with 200000 time iterations and a step-size of  $h = 0.01$ ,  
 280 we numerically solve the system (8). Firstly, we analyze the particular case  
 281 when the coupling strength is  $c = 0$ , the systems are not coupled. Projections  
 282 of the attractors onto the planes  $(x_1^m, x_2^m)$  and  $(x_1^s, x_2^s)$  are given in Figures 3  
 283 a) and c), in both cases the master and slave systems start at the same initial  
 284 condition  $\chi_{m0} = \chi_{s0} = (4.8, 0.48, -0.29)^T$ . This initial condition is indicated

285 with a black dot in figures. The master and slave systems oscillate in a different  
 286 way since they have different scroll degrees  $\eta_m = 3$  and  $\eta_s = 8$ . The elements  
 287 of the index sets  $\mathcal{I}_m = \{8, 9, 10\}$  and  $\mathcal{I}_s = \{3, 4, 5, 6, 7, 8, 9, 10\}$  for the master  
 288 and slave systems, respectively, are indicated on the top of Figures 3 a) and c).  
 289

290 Figures 3 b) and d) show the itineraries  $\mathcal{S}_m(\chi_{m0})$  and  $\mathcal{S}_s(\chi_{s0})$  of the master  
 291 and slave systems, respectively. Note that they are different because the systems  
 292 have different scroll-degrees, even though they start at the same initial condition.  
 293 The itineraries  $\mathcal{S}_m(\chi_{m0})$  and  $\mathcal{S}_s(\chi_{s0})$  are given by the dynamics of the master  
 294 and slave systems and correspond to the activation of the systems in different  
 295 atoms of the partitions, *i.e.*, the itinerary  $\mathcal{S}_m(\chi_{m0})$  generated by  $\kappa_m : \mathbb{R}^3 \rightarrow \mathcal{I}_m$   
 296 only takes three values  $\{8, 9, 10\}$ , meanwhile the itinerary  $\mathcal{S}_s(\chi_{s0})$  is generated  
 by  $\kappa_s : \mathbb{R}^3 \rightarrow \mathcal{I}_s$  and takes eight values  $\{3, 4, 5, 6, 7, 8, 9, 10\}$ .

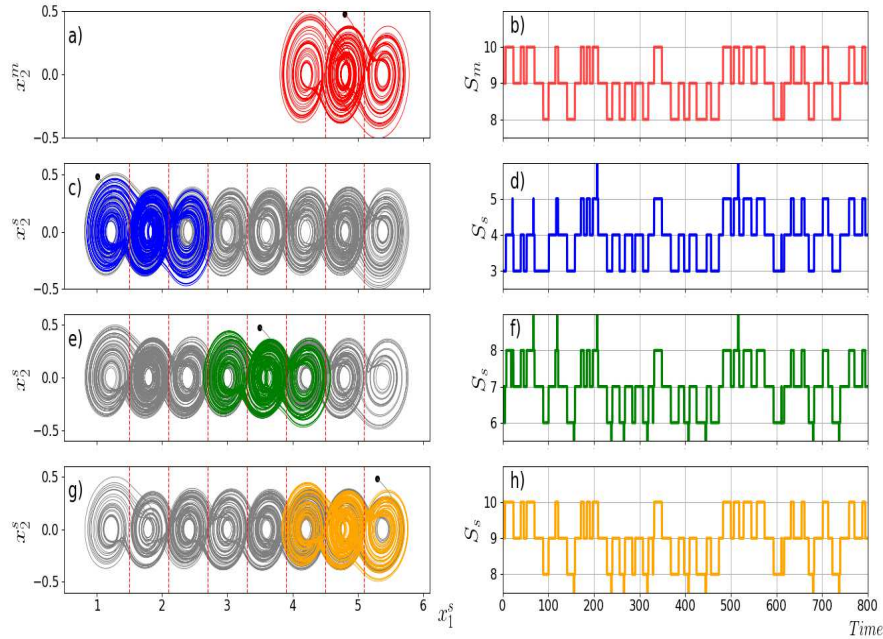


Figure 4: Projections of the master and slave systems onto the  $(x_1^m, x_2^m)$  plane and the  $(x_1^s, x_2^s)$  plane, respectively, for  $\eta_m = 3$ ,  $\eta_s = 8$ ,  $\Gamma = \{0, 1, 0\}$  and coupling strength  $c = 10$ . a) Master system with initial condition  $\chi_{m0} = (4.8, 0.48, -0.29)^T$  and b) its itinerary  $\mathcal{S}_m(\mathcal{X}_{m0})$ . Slave system with different initial conditions: c)  $\chi_{s01} = (1.01, 0.48, -0.29)^T$ , and d) its itinerary  $\mathcal{S}_s(\mathcal{X}_{s01})$ . e)  $\chi_{s02} = (3.5, 0.48, -0.29)^T$  and f) its itinerary  $\mathcal{S}_s(\mathcal{X}_{s02})$ . g)  $\chi_{s03} = (5.3, 0.48, -0.29)^T$  and h) its itinerary  $\mathcal{S}_s(\mathcal{X}_{s03})$ .

297 Now, we set the coupling strength  $c = 10$  and use different initial conditions  
 298 for the slave system. The matrix  $A - c\Gamma$  is Hurwitz for  $0.2 < c$ , with this in mind  
 299 we choose arbitrarily the coupling strength  $c = 10$  to drive the slave system by  
 300 the master system.

301 Figure 4 shows the projections of master-slave system given by (8) onto the  
302 planes  $(x_1^m, x_2^m)$  and  $(x_1^s, x_2^s)$ . Different initial conditions are used for the slave  
303 system located at distinct atoms. For the master system the initial condition is  
304  $\chi_{m0} = (4.8, 0.48, -0.29)^T$ , see Figure 4 a). Specifically we use different initial  
305 conditions for the slave system  $\chi_{s01} = (1.01, 0.48, -0.29)^T$  for Figure 4 c),  $\chi_{s02} =$   
306  $(3.5, 0.48, -0.29)^T$  for Figure 4 e) and  $\chi_{s03} = (5.3, 0.48, -0.29)^T$  for Figure 4 g).  
307 It is worthwhile to observe that the slave system reduces its scroll-degree to  
308 three and, depending on the initial condition, it evolves between distinct basins  
309 of attraction and multistability appears. We plot in gray the trajectory of the  
310 slave system when it is not coupled with the master system in order to compare  
311 it when it is coupled, see Figure 4 c), e) and g).

Notice that the itinerary of the master system  $S_m(\chi_{m0})$  generated by  $\kappa_m : \mathbb{R}^3 \rightarrow \mathcal{I}_m = \{8, 9, 10\}$  remains, however the itinerary of the slave system  $S_m(\chi_{m0})$  generated by  $\kappa_s : \mathbb{R}^3 \rightarrow \mathcal{I}_s$  is determined by its initial condition, for instance, the itinerary takes different values according to the atom where the initial condition belongs  $\chi_{s0} \in P_i$ , for  $i = 3, \dots, 10$ . Now the itinerary of the slave system is restricted to take a subset of the index set  $\mathcal{I}_s$ , *i.e.*,  $\mathcal{I}_s(\chi_{s0}) \subset \mathcal{I}_s$  which will be called restricted index set. This is because the number of scrolls that the slave system coupled with  $c = 10$  displays less scrolls than when it is not coupled. Thus the restricted index sets have different cardinality that is determined by the initial condition  $\chi_{s0} \in P_i$ , for  $i = 3, \dots, 10$ . So for these three initial conditions there are three different restricted index sets given as follows:

$$\kappa_s : \mathbb{R}^3 \rightarrow \mathcal{I}_s(\chi_{s0}) \subset \mathcal{I}_s = \begin{cases} \mathcal{I}_s(\chi_{s01}) = \{3, 4, 5, 6\}, \\ \mathcal{I}_s(\chi_{s02}) = \{5, 6, 7, 8, 9\}, \\ \mathcal{I}_s(\chi_{s03}) = \{7, 8, 9, 10\}. \end{cases}$$

312 The cardinality of the index set  $\mathcal{I}_m$ , and the restricted index sets  $\mathcal{I}_s(\chi_{s01})$ ,  
313  $\mathcal{I}_s(\chi_{s02})$  and  $\mathcal{I}_s(\chi_{s03})$  are 3, 4, 5, and 4, respectively.

314 There is a problem if we want to detect similar behaviour under the presence  
315 of multistability. The inconvenience is resolved by means of defining a new  
316 itinerary based on the trajectory of the systems instead of the dynamics.

Let  $\mathcal{I}_B = \{\#_1, \dots, \#_n\}$  be an index set that labels each element of a partition  $P_\phi = \{P'_1, \dots, P'_n\}$  of the basin of attraction of a dynamical system with flow  $\phi$ . A function  $\kappa : \mathbb{R}^n \rightarrow \mathcal{I}_B$  of the form

$$\kappa(\phi(t, \chi_0)) = \begin{cases} \#_1, & \text{if } \phi(t, \chi_0) \in P'_1; \\ \#_2, & \text{if } \phi(t, \chi_0) \in P'_2; \\ \vdots & \\ \#_n, & \text{if } \phi(t, \chi_0) \in P'_n; \end{cases}$$

317 generates an itinerary of the trajectory. If  $\kappa(\phi(\chi_0)) = s_i \in \mathcal{I}_B$  during the time  
318 interval  $t \in [t_i, t_{i+1})$ , then  $S^\phi(\chi_0) = \{s_0, s_1, s_2, \dots\}$  stands for the itinerary of  
319 the trajectory  $\phi(\chi_0)$ .

320 In our setting in order to describe appropriately the flows of a master-  
321 slave system via symbolic dynamics it is necessary to consider additional atoms  
322  $P_{-n}, \dots, P_0$ , and  $P_{\eta+1}, \dots, P_N$  at the ‘ends’ of the contiguous partition atoms

323 to account for exits and returns to  $P_1$  and  $P_\eta$ , respectively, to the partition  
 324  $\mathcal{P} = \{P_1, \dots, P_\eta\}$ . So we code according to the partition  $\mathcal{P}_\phi = \{P_{-n}, \dots, P_0,$   
 325  $P_1, \dots, P_\eta, P_{\eta+1}, \dots, P_N\}$ . We obtain a symbolic trajectory by writing down the  
 326 sequence of symbols corresponding to the successive partition elements visited  
 327 by the trajectory during a certain period of time.

328 We are interested when the trajectories oscillate in the attractor, so it is  
 329 enough to consider a new partition with two atoms  $P_0$  and  $P_{\eta+1}$  next to the  
 330 atoms  $P_1$  and  $P_\eta$ , *i.e.*,  $\mathcal{P}_\phi = \{P_0, P_1, P_2, \dots, P_\eta, P_{\eta+1}\}$ . So the partition  $\mathcal{P}_\phi$  has  
 331 been obtained by adding two atoms  $P_0$  and  $P_{\eta+1}$  to the partition  $\mathcal{P}$  as follows:

- 332 • The atoms  $P_i \in \mathcal{P}_\phi$ , for  $i = 2, \dots, \eta - 1$ , are the same that the atoms  $P_i \in$   
 333  $\mathcal{P}$ , for  $i = 2, \dots, \eta - 1$ . These atoms are given by the switching surfaces  
 334  $v^T \mathcal{X} = \delta_i$ ,  $i = 1, \dots, \eta - 1$  with  $\delta_2 - \delta_1 = \delta_3 - \delta_2 = \dots = \delta_{\eta-1} - \delta_{\eta-2}$ .
- 335 • The atoms  $P_1, P_\eta \in \mathcal{P}_\phi$  are given by  $P_1 = \{\mathcal{X} \in \mathbb{R}^n : \delta_0 \leq v^T \mathcal{X} < \delta_1\}$ , and  
 336  $P_\eta = \{\mathcal{X} \in \mathbb{R}^n : \delta_{\eta-1} \leq v^T \mathcal{X} < \delta_\eta\}$ , such that  $\delta_1 - \delta_0 = \delta_2 - \delta_1 = \delta_\eta - \delta_{\eta-1}$ .
- 337 • The atoms  $P_0$  and  $P_{\eta+1}$  are given by fulfilling  $\bigcup_{i=0}^{\eta+1} P_i = \mathbb{R}^n$ .

338 For simplicity we generate a new partition  $\mathcal{P}_\phi = \{P_2, P_3, \dots, P_{10}, P_{11}\}$  based  
 339 on the partition  $\mathcal{P} = \{P_3, \dots, P_{10}\}$  which was considered by equation (12),  
 340 because the flow  $\phi(\chi_0) \subset P_\phi$  and the index sets present the same cardinality.

341 The partition  $\mathcal{P}_\phi$  is given as follows:

$$\begin{aligned}
 \mathcal{P}_\phi &= \{ P_2 = \{\mathcal{X} \in \mathbf{R}^3 : x_1 < 0.9\}, \\
 &P_3 = \{\mathcal{X} \in \mathbf{R}^3 : 0.9 \leq x_1 < 1.5\}, P_4 = \{\mathcal{X} \in \mathbf{R}^3 : 1.5 \leq x_1 < 2.1\}, \\
 &P_5 = \{\mathcal{X} \in \mathbf{R}^3 : 2.1 \leq x_1 < 2.7\}, P_6 = \{\mathcal{X} \in \mathbf{R}^3 : 2.7 \leq x_1 < 3.3\}, \\
 &P_7 = \{\mathcal{X} \in \mathbf{R}^3 : 3.3 \leq x_1 < 3.9\}, P_8 = \{\mathcal{X} \in \mathbf{R}^3 : 3.9 \leq x_1 < 4.5\}, \\
 &P_9 = \{\mathcal{X} \in \mathbf{R}^3 : 4.5 \leq x_1 < 5.1\}, P_{10} = \{\mathcal{X} \in \mathbf{R}^3 : 5.1 \leq x_1 < 5.7\}, \\
 &P_{11} = \{\mathcal{X} \in \mathbf{R}^3 : 5.7 \leq x_1\} \}.
 \end{aligned}
 \tag{13}$$

342 Thus  $\mathcal{S}_m^\phi(\mathcal{X}_{m0}) = \{s_0, s_1, \dots, s_m, \dots\}$  stands for the itinerary generated by  
 343 the trajectory of the master system  $\phi_m(t, \mathcal{X}_{m0})$  at  $\mathcal{X}_{m0}$  and,  $\mathcal{S}_m^\phi(i, \mathcal{X}_{m0})$  is the  
 344 element  $s_i \in \mathcal{S}_m^\phi(\mathcal{X}_0)$  that occurs at time  $t_i$ , so the set  $\Delta_{\phi_m} = \{t_0, t_1, \dots, t_m, \dots\}$   
 345 is generated. In a similar way, we can define the itinerary,  $\mathcal{S}_s^\phi(\mathcal{X}_{s0})$  and the set  
 346  $\Delta_{\phi_s} = \{t'_0, t'_1, \dots, t'_m, \dots\}$  generated by the trajectory of the slave system. We  
 347 always assume that the initial conditions belong to their respectively basin of  
 348 attraction of the system.

349 Thereafter, the master index set  $\mathcal{I}_m$  and restricted index sets  $\mathcal{I}_s(\chi_{s01})$ ,  
 350  $\mathcal{I}_s(\chi_{s02})$  and  $\mathcal{I}_s(\chi_{s03})$  have the same cardinality independently of the initial  
 351 conditions  $\chi_{s0} \in P_i$ , for  $i = 3, \dots, 10$ . Now for these three initial conditions  
 352 there are three different restricted index sets with the same cardinality given as  
 353 follows:

$$\kappa_s : \mathbb{R}^3 \rightarrow \mathcal{I}_s(\chi_{s0}) \subset \mathcal{I}_s = \begin{cases} \mathcal{I}_s(\chi_{s01}) = \{2, 3, 4, 5, 6\}, \\ \mathcal{I}_s(\chi_{s02}) = \{5, 6, 7, 8, 9\}, \\ \mathcal{I}_s(\chi_{s03}) = \{7, 8, 9, 10, 11\}. \end{cases}
 \tag{14}$$

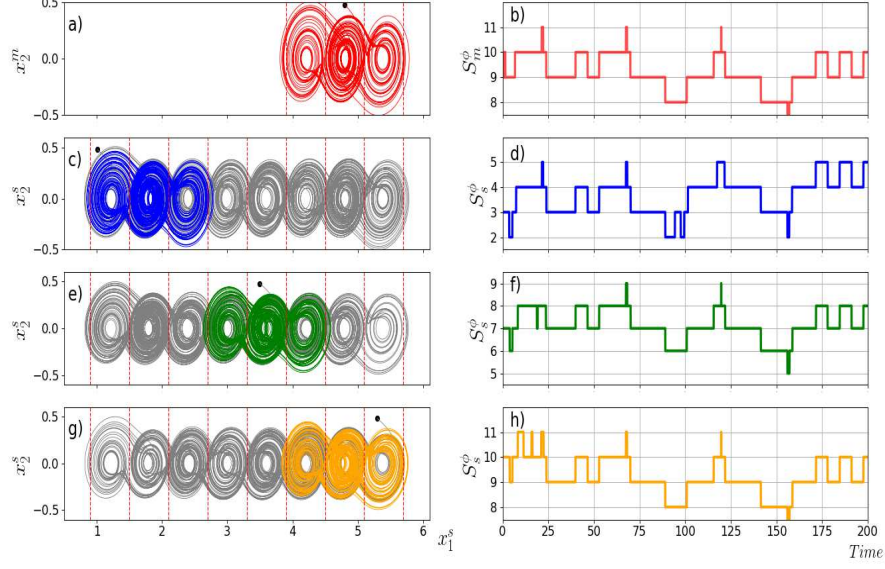


Figure 5: Projections of the master and slave systems onto the  $(x_1^m, x_2^m)$  plane and the  $(x_1^s, x_2^s)$  plane, respectively, for  $\eta_m = 3$ ,  $\eta_s = 8$ ,  $\Gamma = \{0, 1, 0\}$  and coupling strength  $c = 10$ . a) Master system with initial condition  $\chi_{m0} = (4.8, 0.48, -0.29)^T$  and b) its itinerary  $\mathcal{S}_m(\mathcal{X}_{m0})$ . Slave system with different initial conditions: c)  $\chi_{s01} = (1.01, 0.48, -0.29)^T$ , and d) its itinerary  $\mathcal{S}_s(\mathcal{X}_{s01})$ . e)  $\chi_{s02} = (3.5, 0.48, -0.29)^T$  and f) its itinerary  $\mathcal{S}_s(\mathcal{X}_{s02})$ . g)  $\chi_{s03} = (5.3, 0.48, -0.29)^T$  and h) its itinerary  $\mathcal{S}_s(\mathcal{X}_{s03})$ .

And for the master index set:

$$\kappa_m : \mathbb{R}^3 \rightarrow \mathcal{I}_m = \{7, 8, 9, 10, 11\}.$$

354 The cardinality of all of the index set and restricted index sets  $\mathcal{I}_m$ ,  $\mathcal{I}_s(\chi_{s01})$ ,  
355  $\mathcal{I}_s(\chi_{s02})$  and  $\mathcal{I}_s(\chi_{s03})$  is 5. Figure 5 a) shows the projection of the master attractor  
356 onto the plane  $(x_1^m, x_2^m)$  and the atoms of  $P_\phi$  are marked. Figure 5 c), e) and  
357 g) shows the projection of the slave attractor onto the plane  $(x_1^s, x_2^s)$  for different  
358 initial conditions and the atoms of  $P_\phi$  are marked. In Figure 5 b) we show the  
359 itinerary of the master system  $\mathcal{S}_m^\phi(\mathcal{X}_{m0})$  and in Figures 5 d), 5 f) and 5 h) the  
360 itinerary of the slave system by varying the initial condition. Notice that the  
361 itinerary of the trajectory of the master system and the three itineraries of the  
362 trajectories of the slave system for different initial conditions visit five different  
363 domains. Figure 6 shows three signals which were generated by the difference  
364 between the master itinerary  $\mathcal{S}_m^\phi(i, \chi_{m0})$  and slave itineraries for different  
365 initial conditions  $\mathcal{S}_s^\phi(i, \chi_{s0})$ , with  $\chi_{s0} = \{\chi_{s01}, \chi_{s02}, \chi_{s03}\}$ . These signals are comprised  
366 of spikes and a constant offset  $k$ , the spikes correspond to when the trajectory  
367 goes from one atom to other and the constant offset is produced because the  
368 index set  $\mathcal{I}_m$  and restricted index sets  $\mathcal{I}_s(\chi_{s01})$ ,  $\mathcal{I}_s(\chi_{s02})$  and  $\mathcal{I}_s(\chi_{s03})$  are given  
369 by different symbols. The constant offsets  $k_i$ ,  $i = 1, 2, 3$ , by which the average

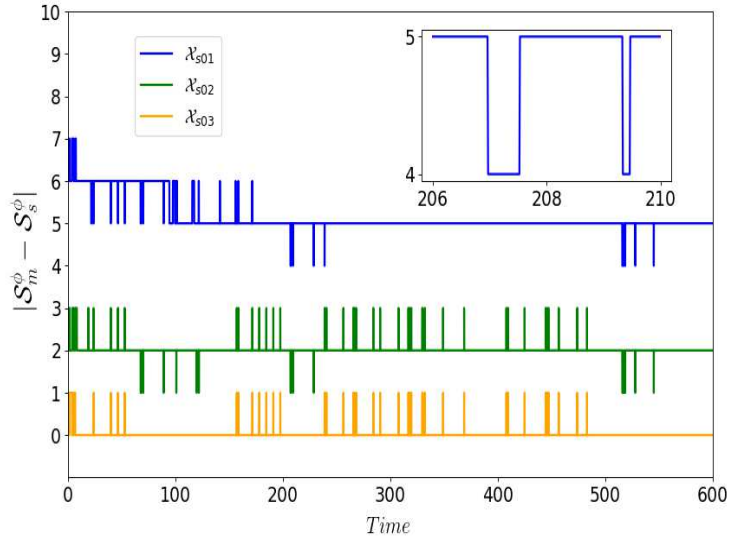


Figure 6: Difference between the itineraries of the master and the slave systems for the initial conditions given in the example 4.3. The inner sub-figure shows a zoom of the blue signal for a short period of time.

370 value of the difference signal is not centered around the  $t$ -axis is computed by  
 371 the  $k_i = |\min\{\mathcal{I}_m\} - \min\{\mathcal{I}_{soi}\}|$ , where  $\min\{\mathcal{I}_j\}$  means the minimum value  
 372 of the set  $\mathcal{I}_j$ . For the initial condition  $\chi_{s01} = (1.01, 0.48, -0.29)^T$  determines  
 373 the constant offset  $k_1 = 5$ ,  $\chi_{s02} = (3.5, 0.48, -0.29)^T$  determines the constant  
 374 offset  $k_2 = 2$  and  $\chi_{s03} = (5.3, 0.48, -0.29)^T$  determines  $k_3 = 0$ . The constant  
 375 offset  $k_3 = 0$  is because the index set  $\mathcal{I}_m$  and the restricted index set  $\mathcal{I}_{s03}$  are  
 376 comprised by the same symbols  $\{7, 8, 9, 10, 11\}$ . If we relabeled the partition  
 377 atoms to make the restricted index sets  $\mathcal{I}_{s01}$  and  $\mathcal{I}_{s02}$  be equal to  $\mathcal{I}_{s03}$ , then all  
 378 the constant offsets  $k_1$ ,  $k_2$  and  $k_3$  will be zero.

## 379 5 Itinerary synchronization

380 In the context of synchronization and multistability, we propose the following  
 381 definition of synchronization based on the itinerary of trajectories in multiscroll  
 382 attractors:

383 *Definition 5.1.* The master-slave system (8) is said to achieve itinerary synchrono-  
 384 zation if after relabeling the partition atoms

$$\lim_{i \rightarrow \infty} |\mathcal{S}_m^\phi(i, \mathcal{X}_{m0}) - \mathcal{S}_s^\phi(i, \mathcal{X}_{s0})| = 0, \quad (15)$$

385 for all initial conditions  $\mathcal{X}_{m0}$  and  $\mathcal{X}_{s0}$  in the basin of attraction.



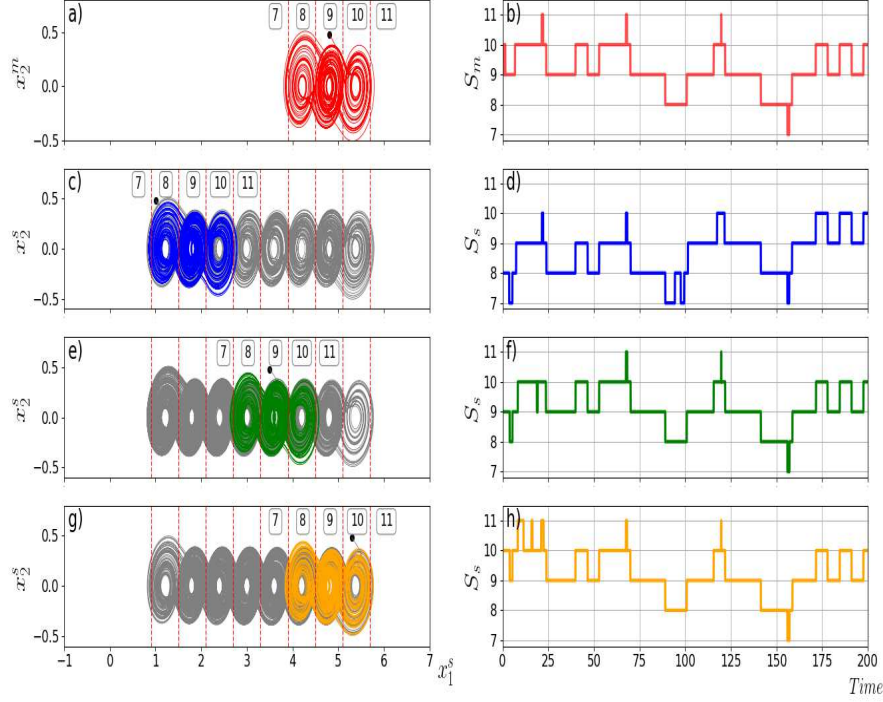


Figure 7: Projections of the master and slave systems onto the  $(x_1^m, x_2^m)$  plane and the  $(x_1^s, x_2^s)$  plane, respectively, for  $\eta_m = 3$ ,  $\eta_s = 8$ ,  $\Gamma = \{0, 1, 0\}$  and coupling strength  $c = 10$ . a) Master system with initial condition  $\chi_{m0} = (4.8, 0.48, -0.29)^T$  and b) its itinerary  $\mathcal{S}_m(\mathcal{X}_{m0})$ . Slave system with different initial conditions: c)  $\chi_{s01} = (1.01, 0.48, -0.29)^T$ , and d) its itinerary  $\mathcal{S}_s(\mathcal{X}_{s01})$  after it was relabeled. e)  $\chi_{s02} = (3.5, 0.48, -0.29)^T$  and f) its itinerary  $\mathcal{S}_s(\mathcal{X}_{s02})$  after it was relabeled. g)  $\chi_{s03} = (5.3, 0.48, -0.29)^T$  and h) its itinerary  $\mathcal{S}_s(\mathcal{X}_{s03})$  without being relabeled.

386 The definition of itinerary synchronization is meant to capture the idea that  
 387 knowing the itinerary of one sequence determines precisely the itinerary of the  
 388 other (after relabeling). Clearly itinerary synchronization will hold if the master-  
 389 slave system (8) presents complete synchronization with the same scroll-degree  
 390 for the master system and slave system since the trajectories of the master and  
 391 slave system will visit the same atoms at the same time.

392 The process of relabeling is shown in Figure 7, here it is possible to see  
 393 that the atoms where the slave system oscillates were relabeled according to the  
 394 master system and we can compare the itineraries between the master and slave  
 395 system. In Figure 7 b) we show the itinerary of the master system  $\mathcal{S}_m^\phi(\mathcal{X}_{m0})$  and  
 396 in Figures 7 d), 7 f) and 7 h) the itinerary of the slave system corresponding to  
 397 various initial conditions,  $\chi_{s01} = (1.01, 0.48, -0.29)^T$ ,  $\chi_{s02} = (3.5, 0.48, -0.29)^T$   
 398 and  $\chi_{s03} = (5.3, 0.48, -0.29)^T$ , respectively.

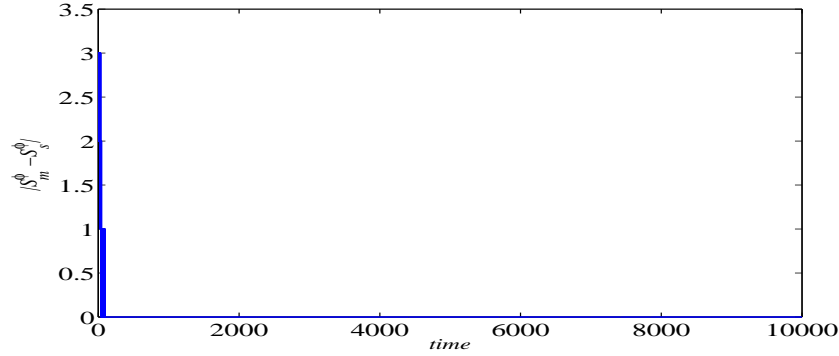


Figure 8: Difference between the itineraries of the master and the slave systems with initial condition:  $\chi_{m0} = (4.8, 0.48, -0.29)^T$  for the master system and  $\chi_{s01} = (1.01, 0.48, -0.29)^T$  for the slave system.

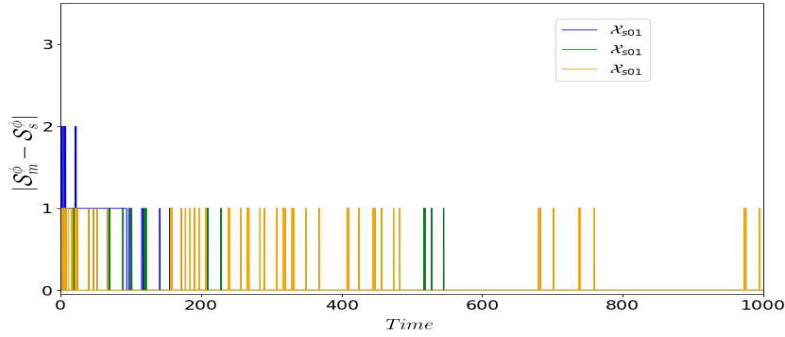


Figure 9: Difference between the itineraries of the master and the slave systems after relabeling the visited atoms for the initial conditions given in the example 4.3.

399 Master and slave systems, the inner linking matrix  $\Gamma$  and the coupling  
400 strength  $c$  play a crucial role to determining whether or not itinerary synchro-  
401 nization holds. For example, if we consider identical systems in the master-slave  
402 system given by (8) and (11) for the master and slave system, with inner linking  
403 matrix  $\Gamma = \text{diag}\{0, 1, 0\}$ , and  $c = 10$ , then itinerary synchronization holds, see  
404 Figure 8. It is worth noting that both systems are identical and oscillate pre-  
405 senting a triple-scroll attractor as shown in Figure 7 a). However, if the systems  
406 are quasi-symmetrical, the master-slave system given by (8), with (11), and (12)  
407 for the master and slave systems respectively, with the same inner linking ma-  
408 trix  $\Gamma$ , and strength coupling  $c$  given previously, itinerary synchronization is lost  
409 for certain recurrent periods of time. Figure (9) shows three signals which were  
410 generated by the difference between the master itinerary and slave itineraries  
411 after relabeling the atoms for different initial conditions for the slave system.

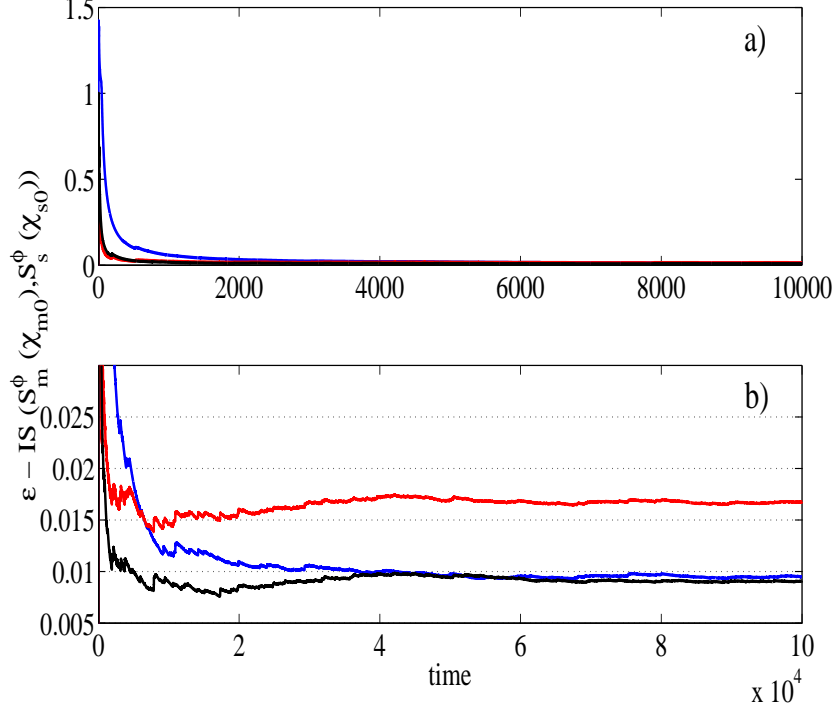


Figure 10: Computation of  $\epsilon$ -Itinerary Synchronization between the master system and slave system after relabeling the visited atoms with the initial condition:  $\chi_{m0} = (4.8, 0.48, -0.29)^T$  for the master system and  $\chi_{s01} = (1.01, 0.48, -0.29)^T$  (blue line),  $\chi_{s02} = (3.5, 0.48, -0.29)^T$  (red line), and  $\chi_{s03} = (5.3, 0.48, -0.29)^T$  (black line) for the slave system, with coupling strength  $c = 10$ . For the time interval of arbitrary units a)  $[0, 10^4]$ ; b)  $[0, 10^5]$ .

412 These small peaks along the error signals indicate that the master and slave  
 413 systems go from one atom to other with an occasional time difference but master  
 414 and slave systems are mostly itinerary synchronized, losing such synchrony  
 415 only when a peaks occurs.

416 The concept of itinerary synchronization is strong for quasi-symmetrical systems.  
 417 A weaker notion of itinerary synchronization is given as follows:

418 *Definition 5.2.* The master-slave system (8) is said to achieve  $\epsilon$ -itinerary syn-  
 419 chronization ( $\epsilon$ -IS) if after relabeling the partition atoms

$$\limsup \frac{1}{t} \int_0^t |S_m^\phi(i, X_{m0}) - S_s^\phi(i, X_{s0})| dt \leq \epsilon \quad (16)$$

420 for all initial conditions  $X_{m0}$  and  $X_{s0}$  in the basin of attraction.

421 The idea of  $\epsilon$ -itinerary synchronization is that the systems are itinerary

422 synchronized except for infrequent (but persistent) time periods. The number  
 423  $\epsilon$  quantifies the asymptotic frequency of asynchronous periods.

424 We investigate  $\epsilon$ -itinerary synchronization of the master-slave system and  
 425 it will be denoted by  $\epsilon - IS(\mathcal{S}_m^\phi, \mathcal{S}_s^\phi)$ . Figure 10 a) shows the computation of  
 426  $\epsilon$ -itinerary synchronization given by (16) when the strength coupling is  $c = 10$   
 427 and the three initial conditions:  $\chi_{s01}$  (blue line),  $\chi_{s02}$  (red line) and  $\chi_{s03}$  (black  
 428 line), previously defined. The master-slave system demonstrates multistability  
 429 and  $\epsilon$ -itinerary synchronization for  $\epsilon = 0.02$ , see Figure 10 b).

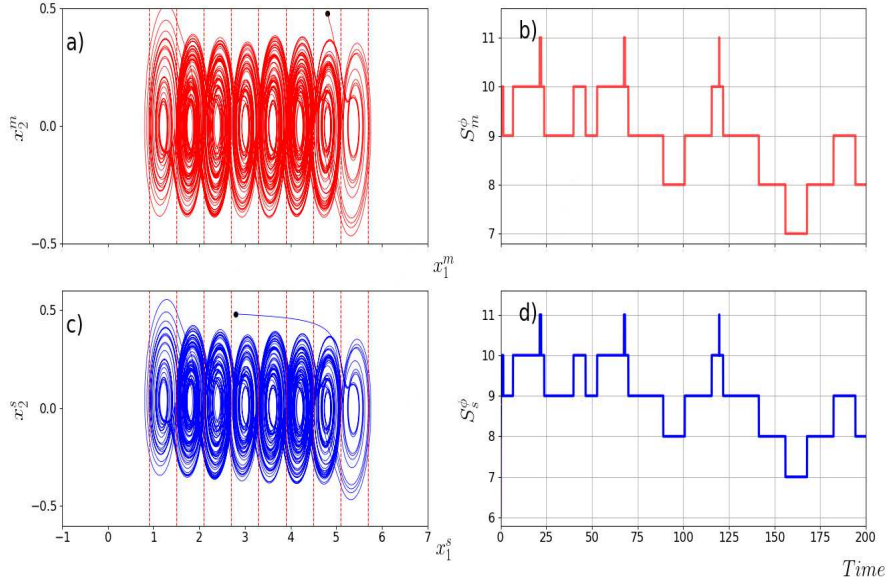


Figure 11: Projections of the master (red color) and slave (blue color) attractors onto the  $(x_1^m, x_2^m)$  plane and the  $(x_1^s, x_2^s)$  plane, respectively, with  $\eta_m = 8$ ,  $\eta_s = 3$ ,  $\Gamma = \{1, 1, 1\}$ , coupling strength  $c = 10$ ; with initial conditions  $\chi_{s_o} = (2.8, 0.48, -0.29)^T$  and  $\chi_{m_o} = (4.8, 0.48, -0.29)^T$  for the slave and master systems, respectively.

430 The multistability phenomenon is given by considering that the scroll degree  
 431 of the master system is less than the scroll-degree of the slave system, and the  
 432 inner linking matrix  $\Gamma = \text{diag}\{0, 1, 0\}$ . By changing the inner linking matrix to  
 433  $\Gamma = \text{diag}\{1, 1, 1\}$  the multistability disappears and the slave system oscillates  
 434 in the same atoms at the same time as the master system as shown in Figure 7 a).  
 435 The inner linking matrix  $\Gamma = \text{diag}\{1, 1, 1\}$  yields the scroll-degree determined  
 436 by the master system in the slave system even if the scroll degree of the master  
 437 system is greater than the scroll-degree of the slave system. For example,  
 438 suppose that the master's scroll-degree is  $\eta_m = 8$  and the slave's scroll-degree  
 439 is  $\eta_s = 3$ . Now the signal for the master system is (12) and for the slave sys-  
 440 tem is (11). We take the inner coupling matrix to be  $\Gamma = \text{diag}\{1, 1, 1\}$  and  
 441 the coupling strength to be  $c = 10$ . This inner coupling matrix  $\Gamma$  makes

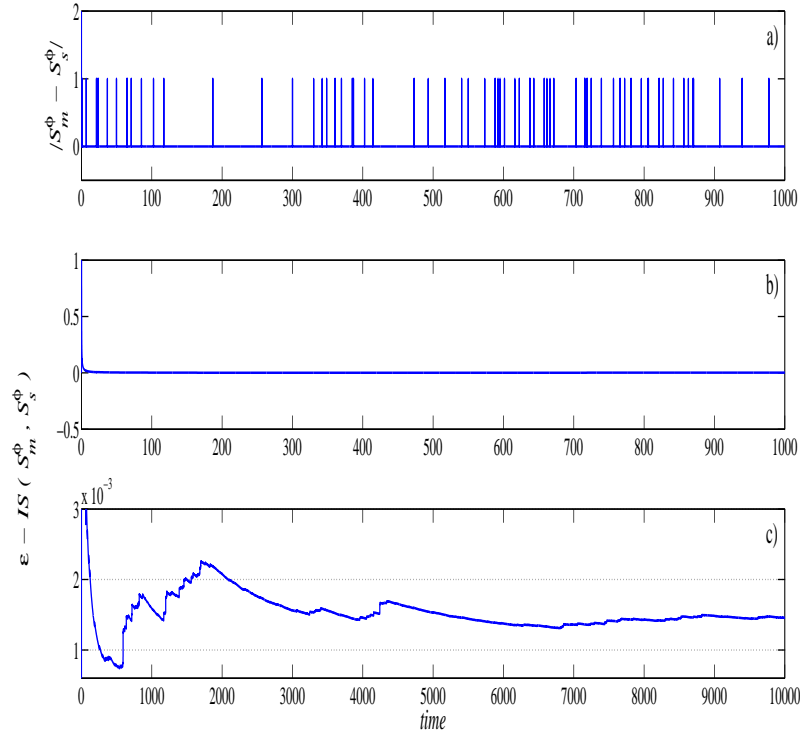


Figure 12: a) shows the difference between the itineraries of the master system and slave system, for  $c = 10$  and  $\Gamma = \text{diag}\{1, 1, 1\}$ . b) and c) show the curve obtained by computing  $\epsilon$ -itineraries of the master system and the slave system. For the time interval of arbitrary units b)  $[0, 10^3]$ ; and c)  $[0, 10^4]$ .

442  $A - c\Gamma$  be Hurwitz for  $0.2 < c$ . Figures 11 a) and c) show the projections of  
443 the master and slave attractors given by (8) onto the  $(x_1^m, x_2^m)$  and  $(x_1^s, x_2^s)$   
444 planes, respectively, generated with initial condition  $\chi_{mo}$  given above for the  
445 master system and  $\chi_{so} = (2.8, 0.48, -0.29)^T$  for the slave system. Note that the  
446 slave system increases its scroll-degree to  $\eta_m = 8$ . Figure 11 b) and d) shows  
447 the master and slave itineraries, respectively. Figure 12 a) shows the differ-  
448 ence between the itineraries of the master system and slave system, for  $c = 10$   
449 and  $\Gamma = \text{diag}\{1, 1, 1\}$ , indicating that the master and slave systems present the  
450 same scroll degree because the offset of the signal is zero. Figure 12 b) shows  
451 the curve obtained by (16) which indicates  $\epsilon$ -itinerary synchronization is achieved  
452 for  $\epsilon = 0.003$ , see Figure 12 c). In this setting the  $\epsilon$  of  $\epsilon$ -Itinerary Synchroniza-  
453 tion tends to zero as the coupling strength increases. This result is shown by  
454 the following proposition 5.3.

455 PROPOSITION (5.3). Consider a master-slave system composed of quasi-sym-  
 456 metrical  $\eta$ -PWL systems described by (8) and signals  $\kappa_m(\mathcal{X})$ , and  $\kappa_s(\mathcal{X})$  given  
 457 by (12) and (11), respectively, with  $\Gamma = \text{diag}\{1, 1, 1\}$  and linear operator  $A$  given  
 458 by (6). As the coupling strength  $c$  tends to infinity then the master-slave system  
 459 presents synchronization.

460 *Proof.* The master slave system is given by

$$\begin{aligned}\dot{\mathcal{X}}_m &= A\mathcal{X}_m + B_{\kappa_m(\mathcal{X}_m)}, \\ \dot{\mathcal{X}}_s &= A\mathcal{X}_s + B_{\kappa_s(\mathcal{X}_s)} + c\Gamma(\mathcal{X}_m - \mathcal{X}_s).\end{aligned}\quad (17)$$

461 Defining the error between the master and slave systems as  $e = \mathcal{X}_m - \mathcal{X}_s =$   
 462  $(e_{x_1}, e_{x_2}, e_{x_3})^T$ , where  $e_{x_1} = x_{m1} - x_{s1}$ ,  $e_{x_2} = x_{m2} - x_{s2}$  and  $e_{x_3} = x_{m3} - x_{s3}$ .  
 463 Thus the error system is given by

$$\begin{aligned}\dot{e} &= Ae + B_{\kappa_m(\mathcal{X}_m)} - B_{\kappa_s(\mathcal{X}_s)} - c\Gamma e, \\ &= (A - c\Gamma)e + B_{\kappa_m(\mathcal{X}_m)} - B_{\kappa_s(\mathcal{X}_s)}, \\ &= \tilde{A}e + B_{\kappa_m(\mathcal{X}_m)} - B_{\kappa_s(\mathcal{X}_s)},\end{aligned}\quad (18)$$

464 So the error system is given by

$$\begin{aligned}\dot{e}_{x_1} &= -ce_{x_1} + e_{x_2}, \\ \dot{e}_{x_2} &= -ce_{x_2} + e_{x_3}, \\ \dot{e}_{x_3} &= -\alpha_{31}e_{x_1} - \alpha_{32}e_{x_2} - (\alpha_{33} + c)e_{x_3} - (\beta_m - \beta_s),\end{aligned}\quad (19)$$

where  $\beta_m$  and  $\beta_s$  take values of the third entry of the vectors  $B_j$ , with  $j =$   
 3, 4, 5, 6, 7, 8, 9, 10 and  $B_j$ , with  $j = 8, 9, 10$ , respectively. Solving for the equi-  
 librium point we find

$$e_{x_1} = (\beta_m - \beta_s)/(-\alpha_{31} - \alpha_{32} - (\alpha_{33} + c)c^2).$$

465 As  $(\beta_m - \beta_s)$  is bounded  $e_{x_1}$  tends to zero when  $c$  tends to infinity. If  $e_{x_1}$   
 466 tends to zero then  $e_{x_2}$  and  $e_{x_3}$  also tend to zero. Therefore, the error system  
 467 has  $(0, 0, 0)^T$  as its sole equilibrium point. The master-slave system displays  
 468 synchronization.  $\square$

469 In our numerical results we have considered only  $c = 10$  and  $\Gamma = \text{diag}\{0, 1, 0\}$   
 470 and  $\Gamma = \text{diag}\{1, 1, 1\}$ . But for sufficiently large values of  $c$  with both  $\Gamma =$   
 471  $\text{diag}\{0, 1, 0\}$  and  $\Gamma = \text{diag}\{1, 1, 1\}$  the matrix  $A - c\Gamma$  will have only eigenval-  
 472 ues with negative real part, which should lead to itinerary synchronization or  
 473  $\epsilon$ -itinerary synchronization for small  $\epsilon$  tending to 0 as the coupling strength  
 474 increases (the presence of discontinuities in the PWL system makes difficult a  
 475 rigorous rather than heuristic proof). For example if  $\Gamma = \text{diag}\{1, 1, 1\}$  and  $c$  is  
 476 greater than the positive real part of conjugate eigenvalues  $\lambda_2$  and  $\lambda_3$  of  $A$  then  
 477  $A - c\Gamma$  will have all eigenvalues with negative real part.

## 478 6 Dynamical Networks

479 A dynamical network is composed of  $N$  coupled dynamical systems called nodes  
 480 [20]. Each node is labeled by an index  $i = 1, \dots, N$  and described by a first  
 481 ordinary differential equation system of the form  $\dot{\mathcal{X}}_i(t) = f_i(\mathcal{X}_i(t))$ , where  
 482  $\mathcal{X}_i(t) = (x_{i1}(t), \dots, x_{in}(t))^T \in \mathbb{R}^n$  is the state vector and,  $f_i : \mathbb{R}^n \rightarrow \mathbb{R}^n$  is  
 483 the vector field which describes the dynamical behavior of an  $i$ -th node when it  
 484 is not connected to the network. We assume the coupling between neighboring  
 485 nodes is linear so that the state equation of the entire network is described by  
 486 the following equations:

$$\dot{\mathcal{X}}_i(t) = f_i(\mathcal{X}_i(t)) + c \sum_{j=1}^N \Delta_{ij} \Gamma (\mathcal{X}_j(t) - \mathcal{X}_i(t)), \quad i = 1, \dots, N, \quad (20)$$

487 where  $c$  is the uniform coupling strength between the nodes and the inner  
 488 linking matrix  $\Gamma = \text{diag}\{r_1, \dots, r_n\} \in \mathbb{R}^{n \times n}$  is described in (8). The matrix  
 489  $\Delta = \{\Delta_{ij}\} \in \mathbb{R}^{N \times N}$  is called a coupling matrix if its elements are zero or  
 490 one depending on which nodes are connected or not. Such a matrix contains  
 491 the entire information about the network configuration topology. Specifically, if  
 492 nodes are coupled with bidirectional links, then  $\Delta$  is a symmetric matrix with  
 493 the following entries: if there is a connection between node  $i$  and node  $j$  (with  
 494  $i \neq j$ ), then  $\Delta_{ij} = \Delta_{ji} = 1$ ; otherwise  $\Delta_{ij} = \Delta_{ji} = 0$ .

495 On the other hand, if the nodes are connected with unidirectional links,  
 496 then  $\Delta$  is a non-symmetric matrix with entries defined as follows:  $\Delta_{ij} = 1$   
 497 (with  $i \neq j$ ) if there is an edge directed from node  $j$  to node  $i$ ;  $\Delta_{ij} = 0$  if node  
 498  $j$  is not connected to node  $i$ .

499 Network (20) can be equivalently expressed in matrix form by using the  
 500 Kronecker product as follows:

$$\dot{X}(t) = F(X(t)) + c(\Delta \otimes \Gamma)X(t),$$

501 where  $X(t) = (\mathcal{X}_1, \dots, \mathcal{X}_N)^T \in \mathbb{R}^{Nn}$ ;  $F(X(t)) = (f_1(\mathcal{X}_1), \dots, f_N(\mathcal{X}_N))^T \in$   
 502  $\mathbb{R}^{Nn}$ ; and  $\otimes$  denotes the Kronecker product of matrices.

503 For the dynamical network (20) with a symmetric coupling matrix, one of  
 504 the most studied collective phenomena is synchronization, which emerges when  
 505 the dynamical behavior between nodes are correlated in-time (See [20] and refer-  
 506 ences there in).

## 507 7 Ring and chain topology networks

508 We study the collective dynamics of  $N$  coupled quasi-symmetrical  $\eta$ -PWL sys-  
 509 tems which are connected by unidirectional links in a ring topology, *i.e.*, a  
 510 network composed of an ensemble of master-slave systems coupled in a cascade  
 511 configuration topology. In this context, a system defined in the node  $i$  is a slave  
 512 system of a system defined in the node  $i - 1$ , and also plays the role of a master  
 513 system for a system defined in the node  $i + 1$ . Figure 13 (a) shows a network

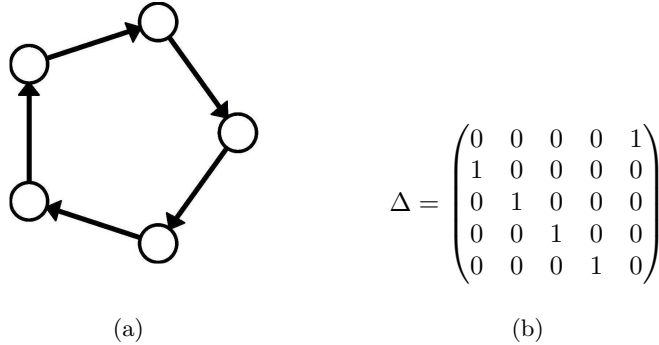


Figure 13: A network of  $N = 5$  nodes coupled in a ring topology with unidirectional links. a) The network topology and b) the coupling matrix.

514 with a ring topology and 13 (b) its corresponding coupling matrix  $\Delta$ . A network  
 515 with such attributes is described by the following state equations:

$$\begin{cases} \dot{\mathcal{X}}_1 = A\mathcal{X}_1 + B_{\kappa_1}(\mathcal{X}_1) + c\Gamma(\mathcal{X}_N - \mathcal{X}_1), \\ \dot{\mathcal{X}}_2 = A\mathcal{X}_2 + B_{\kappa_2}(\mathcal{X}_2) + c\Gamma(\mathcal{X}_1 - \mathcal{X}_2), \\ \dot{\mathcal{X}}_3 = A\mathcal{X}_3 + B_{\kappa_3}(\mathcal{X}_3) + c\Gamma(\mathcal{X}_2 - \mathcal{X}_3), \\ \vdots \\ \dot{\mathcal{X}}_N = A\mathcal{X}_N + B_{\kappa_N}(\mathcal{X}_N) + c\Gamma(\mathcal{X}_{N-1} - \mathcal{X}_N), \end{cases} \quad (21)$$

516 where  $\mathcal{X}_i$ ,  $i = 1, 2, \dots, N$ , denotes the state vector of each node. Notice that the  
 517 system (21) is a dynamical network where each node differs only in the constant  
 518 vector  $B_{\kappa_i}(\cdot)$ . In this context, we propose the following definition of a network  
 519 of nearly identical nodes:

520 *Definition 7.1.* A network of nearly identical nodes is a network composed of  
 521 nodes with dynamics given by quasi-symmetrical  $\eta$ -PWL systems, *i.e.*,  $A_i =$   
 522  $A_j = A$ ,  $\eta_i \neq \eta_j$  and  $\kappa_i(\cdot) \neq \kappa_j(\cdot) \forall i, j = 1, 2, \dots, N$  whose state equation is  
 523 written as follows:

$$\dot{\mathcal{X}}_i = A\mathcal{X}_i + B_{\kappa_i}(\mathcal{X}_i) + c \sum_{j=1}^N \Delta_{ij} \Gamma(\mathcal{X}_j - \mathcal{X}_i), \quad i = 1, \dots, N. \quad (22)$$

524 Note that (22) corresponds to a dynamical network with a configuration  
 525 topology given by the coupling matrix  $\Delta = \{\Delta_{ij}\} \in \mathbb{R}^{N \times N}$ . In particular, for a  
 526 ring topology (Figure 13), the equation (22) becomes the equation (21).

527 We first study the collective behavior of a nearly identical network (22)  
 528 assuming that the coupling matrix corresponds to a network with a ring topology  
 529 and with unidirectional links. We are interested in knowing what is the scroll-  
 530 degree of all the nodes in this kind of network with different scroll-degree in its  
 531 nodes and when none of them is the leading node (master system). In Section 5,



532 a master system forces the slave system to have the same scroll degree and the  
533 master-slave system achieves  $\epsilon$ -Itinerary Synchronization. However in a ring  
534 topology network each node behaves as the master system of the following node  
535 but at the same time it behaves as a slave system of the preceding node. Finally  
536 we consider the case in which the network (22) has a directed chain topology  
537 where the leading node has the maximum or the minimum scroll-degree.

## 538 7.1 Node's dynamics

539 In this section we consider the switching regions as in Definition 2.1. The  
540 dynamics of the  $i$ -node is controlled by the  $(i-1)$ -node, see equation (21).

541 Since the dynamics of a single node is governed by an UDS system plus a  
542 coupling signal which comes from only one node of the network, we know that  
543 the linear operator is diagonalizable *i.e.* exist a matrix  $Q \in \mathbb{R}^{3 \times 3}$  such that  
544  $\Lambda = Q^{-1}AQ$  with  $\Lambda = \text{diag}[\lambda_1, \lambda_2, \lambda_3]$ . So the node's dynamics is given by

$$\begin{aligned} \dot{x}_{i1} &= -cx_{i1} + x_{i2} + cx_{(i-1)1}, \\ \dot{x}_{i2} &= -cx_{i2} + x_{i3} + cx_{(i-1)2}, \\ \dot{x}_{i3} &= -1.5x_{i1} - x_{i2} - (1+c)x_{i3} + \beta_3^{\kappa_i} + cx_{(i-1)3}, \end{aligned} \quad (23)$$

545 where  $\mathcal{X}_i = (x_{i1}, x_{i2}, x_{i3})^T$ , for  $i = 1, \dots, N$  and consider that if  $n = 1$  then  
546  $n - 1 = N$ .  $\beta_3^{\kappa_i} \in \Delta_\beta = \{0, 0.9, 1.8, 2.7, 3.6, 4.5, 5.4, 6.3, 7.2, 8.1\}$  is determined  
547 by the third component of constant vectors  $B_{\kappa_i} = (0, 0, \beta_3^{\kappa_i})$ . By introducing a  
548 change of variable  $z_i = (z_1^i, z_2^i, z_3^i)^T = (x_{i1} - k_1, x_{i2} - k_2, x_{i3} - k_3)^T$  each the  
549 trajectory  $\mathcal{X}_i(t)$  goes to an atom  $P_i$  of the partition  $\mathcal{P}$ , with  $k_1 = \beta_3^{\kappa_i}/(1.5 + c +$   
550  $c^2 + c^3)$ ,  $k_2 = cK_1$ , and  $k_3 = c^2k_1$ . We rewrite the equation (23) as follows:

$$\begin{aligned} \dot{z}_1^i &= -cz_1^i + z_2^i + f_1, \\ \dot{z}_2^i &= -cz_2^i + z_3^i + f_2, \\ \dot{z}_3^i &= -1.5z_1^i - z_2^i - (1+c)z_3^i + f_3, \end{aligned} \quad (24)$$

551 where  $f_1 = cx_{(i-1)1}$ ,  $f_2 = cx_{(i-1)2}$ , and  $f_3 = cx_{(i-1)3}$  are external signal of the  
552  $i$ -node that come from  $(i-1)$ -node. So the system (24) is given as follows:

$$\dot{z}_i = A_c z_i + F^{i-1}, \quad i = 1, \dots, N, \quad (25)$$

553 where  $A_c = A + \text{diag}[-c, -c, -c]$  and  $F^{i-1} = [f_1, f_2, f_3]^T$  is conformed from  
554 the state vector of the  $(i-1)$ -node. If  $c > 0.1020$  then  $A_c$  is Hurwitz. For  
555 the particular value of  $c = 10$  the eigenvalues are:  $\lambda_{c1} = -11.2041$ ,  $\lambda_{c2} =$   
556  $-9.8980 + 1.1115i$ ,  $\lambda_{c3} = -9.8980 - 1.1115i$ . The solution of the nonhomogeneous  
557 linear system (25) is:

$$z_i(t) = e^{A_c t} z_i(0) + e^{A_c t} \int_0^t e^{-A_c \tau} F(\tau) d\tau, \quad (26)$$

558 where  $z_i(0)$  is the initial condition of the  $i$ -th node in the new state variable.  
559 The first term of the right hand side of the equation (26) converges to zero when

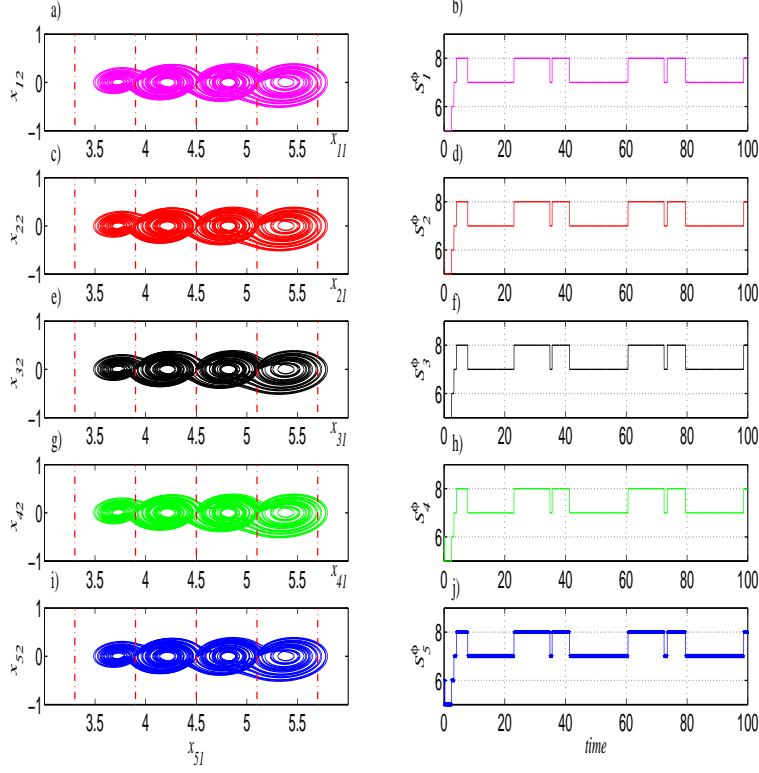


Figure 14: Dynamics of a nearly identical network (22) with coupling strength  $c = 10$  and  $\Gamma = \text{diag}\{1, 1, 1\}$ ; the scroll-degree and initial condition for each node are given in Table (1): a); c); e); g); i); The projections of the attractors onto the plane  $(x_{i1}, x_{i2})$  of the node 1,2,3,4 and 5 respectively (Transient were removed); and b); d); f); h); j) its itinerary.

560  $t \rightarrow \infty$ . So the node's dynamics is given as follows

$$\mathcal{X}_i(t) = (k_1, k_2, k_3)^T + e^{A_c t} \int_0^t e^{-A_c \tau} \mathcal{X}_{i-1}(\tau) d\tau. \quad (27)$$

561 The dynamics of  $i$ -node is determined by the  $(i - 1)$ -node, so the collective  
 562 dynamics of  $N$  coupled quasi-symmetrical  $\eta$ -PWL systems which are connected  
 563 by unidirectional links in a ring topology can present synchronous behavior if  
 564 the different node states commute from one atom  $P_i$  to other  $P_j$  presenting the  
 565 same constant vector  $(k_1, k_2, k_3)^T$ .

566 **7.2 Dynamics in a ring topology**

567 We consider a ring network with five nodes, *i.e.*,  $N = 5$  nearly identical nod-  
 568 es described in (22) and coupled in a ring topology. We assume that each  
 569 node's dynamic is described by the same linear operator  $A$  (*i.e.* they are quasi-  
 570 symmetrical) and a subset of the set of constant vectors  $\mathbf{B} = \{B_1, B_2, \dots, B_{10}\}$   
 571 which are those given by (6). Further, for each node we select the scroll-degree  
 572 ( $\eta_i$ ) and its corresponding initial condition according to Table (1).

Node's label	Scroll-degree	Initial condition
1	10	$(0.227, -0.216, -0.359)^T$
2	5	$(3.014, -0.371, -0.271)^T$
3	3	$(5.349, -0.424, -0.279)^T$
4	8	$(1.402, -0.205, -0.316)^T$
5	6	$(2.452, -0.266, -0.308)^T$

Table 1: The scroll-degree ( $\eta_i$ ) and its corresponding initial condition for each node in the nearly identical network coupled in a ring topology for Examples 7.2 and 7.3.

573 The signal for the first node with scroll-degree  $\eta_1 = 10$  is given by (7) where  
 574 is defined the partition  $\mathcal{P} = \{P_1, \dots, P_{10}\}$ ; for the third and fourth nodes with  
 575 scroll-degree  $\eta_3 = 3$  and partition  $\mathcal{P} = \{P_8, \dots, P_{10}\}$ ; and  $\eta_4 = 8$  and partition  
 576  $\mathcal{P} = \{P_3, \dots, P_{10}\}$  are given by (11) and (12) respectively. For the second node  
 577 with scroll degree  $\eta_2 = 5$  the switching signal is given as follows:

$$\kappa_5(\mathcal{X}) = \begin{cases} 1, & \text{if } \mathcal{X} \in P_{10} = \{\mathcal{X} \in \mathbb{R}^3 : x_1 \geq 5.1\}; \\ 2, & \text{if } \mathcal{X} \in P_9 = \{\mathcal{X} \in \mathbb{R}^3 : 4.5 \leq x_1 < 5.1\}; \\ 3, & \text{if } \mathcal{X} \in P_8 = \{\mathcal{X} \in \mathbb{R}^3 : 3.9 \leq x_1 < 4.5\}; \\ 4, & \text{if } \mathcal{X} \in P_7 = \{\mathcal{X} \in \mathbb{R}^3 : 3.3 \leq x_1 < 3.9\}; \\ 5, & \text{if } \mathcal{X} \in P_6 = \{\mathcal{X} \in \mathbb{R}^3 : x_1 < 3.3\}. \end{cases} \quad (28)$$

578 And for the fifth node with scroll degree  $\eta_5 = 6$  is

$$\kappa_6(\mathcal{X}) = \begin{cases} 1, & \text{if } \mathcal{X} \in P_{10} = \{\mathcal{X} \in \mathbb{R}^3 : x_1 \geq 5.1\}; \\ 2, & \text{if } \mathcal{X} \in P_9 = \{\mathcal{X} \in \mathbb{R}^3 : 4.5 \leq x_1 < 5.1\}; \\ 3, & \text{if } \mathcal{X} \in P_8 = \{\mathcal{X} \in \mathbb{R}^3 : 3.9 \leq x_1 < 4.5\}; \\ 4, & \text{if } \mathcal{X} \in P_7 = \{\mathcal{X} \in \mathbb{R}^3 : 3.3 \leq x_1 < 3.9\}; \\ 5, & \text{if } \mathcal{X} \in P_6 = \{\mathcal{X} \in \mathbb{R}^3 : 2.7 \leq x_1 < 3.3\}; \\ 6, & \text{if } \mathcal{X} \in P_5 = \{\mathcal{X} \in \mathbb{R}^3 : x_1 < 2.7\}. \end{cases} \quad (29)$$

579 The scroll-degree is determined numerically under two inner coupling ma-  
 580 trices:  $\Gamma = \text{diag}\{1, 1, 1\}$  and  $\Gamma = \text{diag}\{1, 0, 0\}$ , and the ring topology network  
 581 with five nodes.

582 **Example 7.2.** For the nearly identical network described above, we assume  
 583 that the coupling strength is  $c = 10$ , the inner coupling matrix is  $\Gamma = \text{diag}\{1, 1, 1\}$ .  
 584 We solve numerically the nearly identical network (22) with the scroll-degree

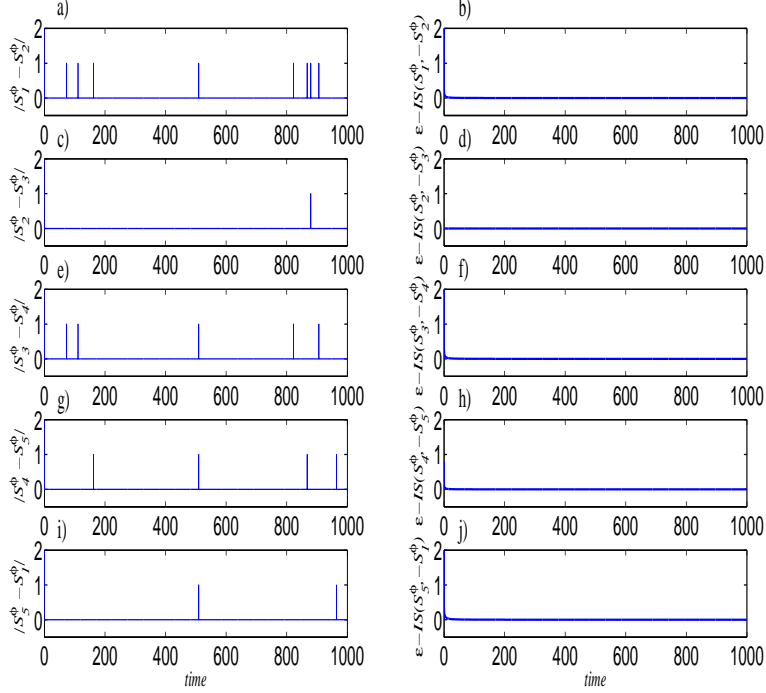


Figure 15: Difference between the itineraries of the nodes of a nearly identical network (22) with coupling strength  $c = 10$  and  $\Gamma = \text{diag}\{1, 1, 1\}$ ; the scroll-degree and initial condition for each node are given in Table (1). a)  $|\mathcal{S}_1^\phi - \mathcal{S}_2^\phi|$ ; c)  $|\mathcal{S}_2^\phi - \mathcal{S}_3^\phi|$ ; e)  $|\mathcal{S}_3^\phi - \mathcal{S}_4^\phi|$ ; g)  $|\mathcal{S}_4^\phi - \mathcal{S}_5^\phi|$ ; i)  $|\mathcal{S}_5^\phi - \mathcal{S}_1^\phi|$ ; and b)  $\epsilon - IS(\mathcal{S}_1^\phi, \mathcal{S}_2^\phi)$ ; d)  $\epsilon - IS(\mathcal{S}_2^\phi, \mathcal{S}_3^\phi)$ ; f)  $\epsilon - IS(\mathcal{S}_3^\phi, \mathcal{S}_4^\phi)$ ; h)  $\epsilon - IS(\mathcal{S}_4^\phi, \mathcal{S}_5^\phi)$ ; j)  $\epsilon - IS(\mathcal{S}_5^\phi, \mathcal{S}_1^\phi)$ .

585 and initial condition given in Table (1) and using a Runge-Kutta method with  
 586 10,000,000 time iterations and step size  $h = 0.01$ .

587 In the first column of the Figure 14 we show the projections of the attractors  
 588 onto the planes  $(x_{i1}, x_{i2})$  after transients, with  $i = 1, \dots, 5$ , note that  
 589 independently of the initial conditions, the trajectories of all nodes converge to  
 590 an attractor with four scrolls and one of them is a smaller scroll than the others  
 591 (the left scroll). If we count this smaller scroll, then the ring topology network  
 592 displays a  $\eta = 4$  scroll degree. In the right column of the Figure 14, we display  
 593 its corresponding itinerary in a short interval of time in order to appreciate the  
 594 time elapsed that the trajectory of each node spends in a given atom. We can  
 595 see that in this short time the itineraries behave identically and definition of  
 596 itinerary synchronization is fulfilled. However if we analyze the difference be-  
 597 tween itineraries of the  $(i-1)$ -th node and  $i$ -th node in a longer period of time  
 598 it is possible to see that the nodes are briefly out of itinerary synchronization.  
 599 For example, Figure 15 a) shows the difference of itineraries of the first node

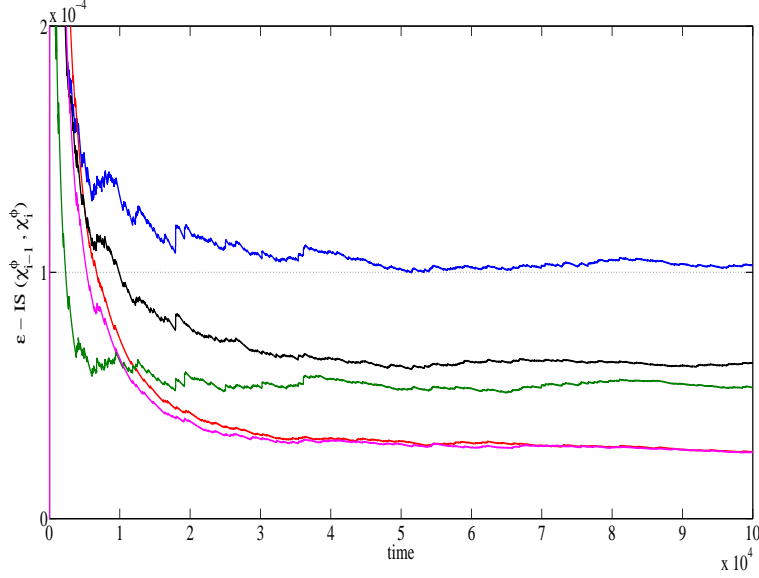


Figure 16: Different curves computed by (16) for  $\epsilon$ -itinerary synchronization between the nodes of a nearly identical network (22) with coupling strength  $c = 10$  and  $\Gamma = \text{diag}\{1, 1, 1\}$ ; the scroll-degree and initial condition for each node are given in Table (1): blue line for  $\epsilon - IS(\mathcal{S}_1^\phi, \mathcal{S}_2^\phi)$ ; red line for  $\epsilon - IS(\mathcal{S}_2^\phi, \mathcal{S}_3^\phi)$ ; black line for  $\epsilon - IS(\mathcal{S}_3^\phi, \mathcal{S}_4^\phi)$ ; green line for  $\epsilon - IS(\mathcal{S}_4^\phi, \mathcal{S}_5^\phi)$ ; and magenta line for  $\epsilon - IS(\mathcal{S}_5^\phi, \mathcal{S}_1^\phi)$ .

600 and the second node  $|\mathcal{S}_1^\phi - \mathcal{S}_2^\phi|$ , remember that the coupling is unidirectional,  
601 *i.e.*, the first node acts as a master system on the second node which acts as  
602 a slave system. These two nodes are synchronized when the difference between  
603 itineraries is zero and out of synchronization otherwise. Figure 15 shows the dif-  
604 ference of itineraries of: c) the second node and the third node  $|\mathcal{S}_2^\phi - \mathcal{S}_3^\phi|$ ; e) the  
605 third node and the fourth node  $|\mathcal{S}_3^\phi - \mathcal{S}_4^\phi|$ ; g) the fourth node and the fifth node  
606  $|\mathcal{S}_4^\phi - \mathcal{S}_5^\phi|$ ; and i) the fifth node and the first node  $|\mathcal{S}_5^\phi - \mathcal{S}_1^\phi|$ . Figure 15 b) shows  
607 the  $\epsilon$ -itinerary synchronization between the first node and the second node, it is  
608 possible to see that  $\epsilon$ -itinerary synchronization definition is fulfilled. Figures 15  
609 d), f), h), and j) show the  $\epsilon$ -itinerary synchronizations between the  $i$ -th node  
610 and its  $(i+1)$ -th node: d)  $\epsilon - IS(\mathcal{S}_2^\phi, \mathcal{S}_3^\phi)$ ; f)  $\epsilon - IS(\mathcal{S}_3^\phi, \mathcal{S}_4^\phi)$ ; h)  $\epsilon - IS(\mathcal{S}_4^\phi, \mathcal{S}_5^\phi)$ ;  
611 j)  $\epsilon - IS(\mathcal{S}_5^\phi, \mathcal{S}_1^\phi)$ . In conclusion, all the nodes of the ring topology network  
612 present  $\epsilon$ -itinerary synchronization by considering  $\epsilon = 0.0002$ , see Figure 16.  
613 This Figure shows the different curves computed by (16) for  $\epsilon$ -itinerary syn-  
614 chronization between the nodes of a nearly identical network (22) with coupling  
615 strength  $c = 10$  and  $\Gamma = \text{diag}\{1, 1, 1\}$ ; the scroll-degree and initial condition  
616 for each node are given in Table (1): blue line for  $\epsilon - IS(\mathcal{S}_1^\phi, \mathcal{S}_2^\phi)$ ; red line for  
617  $\epsilon - IS(\mathcal{S}_2^\phi, \mathcal{S}_3^\phi)$ ; black line for  $\epsilon - IS(\mathcal{S}_3^\phi, \mathcal{S}_4^\phi)$ ; green line for  $\epsilon - IS(\mathcal{S}_4^\phi, \mathcal{S}_5^\phi)$ ; and  
618 magenta line for  $\epsilon - IS(\mathcal{S}_5^\phi, \mathcal{S}_1^\phi)$ .

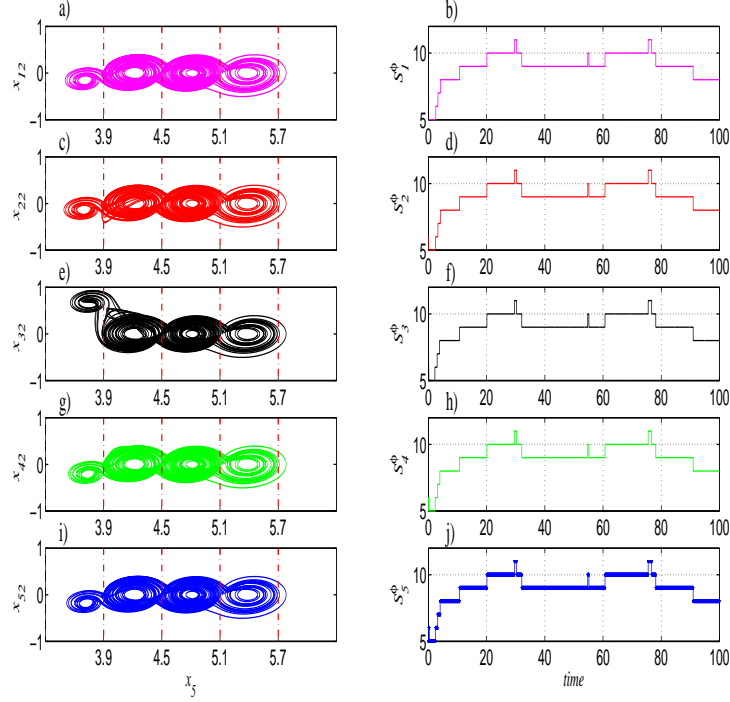


Figure 17: Dynamics of a nearly identical network (22) with coupling strength  $c = 10$  and  $\Gamma = \text{diag}\{1, 0, 0\}$ ; the scroll-degree and initial condition for each node are given in Table (1): a); c); e); g); i); The projections of the attractors onto the plane  $(x_1, x_2)$  of the node 1,2,3,4 and 5 respectively (Transients were removed); and b); d); f); h); j) its itinerary.

619 **Example 7.3.** The dynamics of the network composed of  $N$  quasi-symmetrical  
620  $\eta$ -PWL systems described above can display several behaviors depending on the  
621 inner coupling matrix  $\Gamma$ . The collective dynamics is affected when we suppress  
622 some variable state in the inner connection. For example, in the first column of  
623 Figure 17 when we suppress two state variables from the inner coupling matrix  
624  $\Gamma = \text{diag}\{1, 0, 0\}$ , a deformation of the scroll attractor is achieved specially over  
625 the node with the smallest node-degree (in this case, for the node with scroll-  
626 degree 3). In the first column of the Figure 17 we show the projections of the  
627 attractors onto the planes  $(x_{i1}, x_{i2})$ , with  $i = 1, \dots, 5$ , the ring topology network  
628 displays a  $\eta = 4$  scroll degree. In the right column of the Figure 17, we display  
629 its corresponding itinerary in a short interval of time in order to appreciate the  
630 time elapsed that the trajectory of each node spends in a given atom. We can  
631 see that in this short time the itineraries behave identically and definition of  
632 itinerary synchronization is fulfilled again that for  $\Gamma = \text{diag}\{1, 1, 1\}$ . And the  
633 difference between itineraries of the  $(i - 1)$ -th node and  $i$ -th node is shown

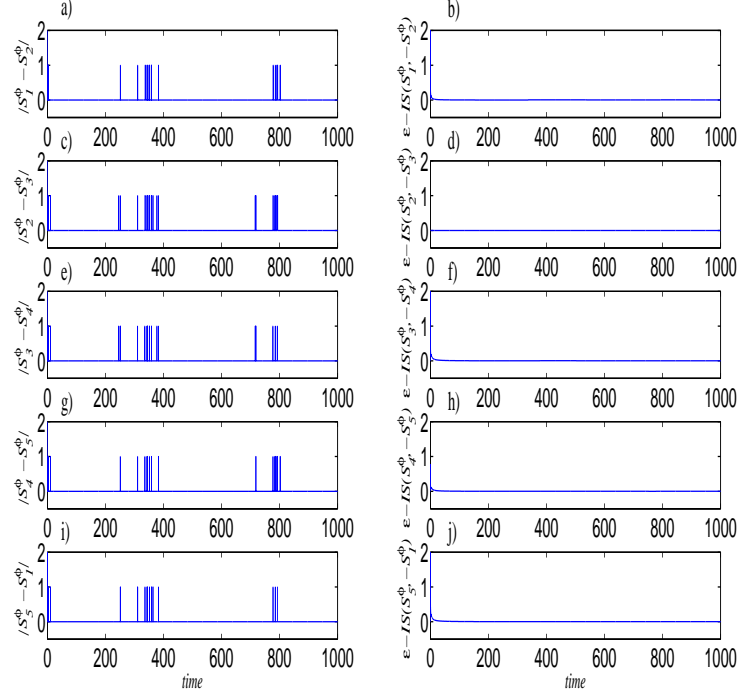


Figure 18: Difference between the itineraries of the nodes of a nearly identical network (22) with coupling strength  $c = 10$  and  $\Gamma = \text{diag}\{1, 0, 0\}$ ; the scroll-degree and initial condition for each node are given in Table (1). a)  $|\mathcal{S}_1^\phi - \mathcal{S}_2^\phi|$ ; c)  $|\mathcal{S}_2^\phi - \mathcal{S}_3^\phi|$ ; e)  $|\mathcal{S}_3^\phi - \mathcal{S}_4^\phi|$ ; g)  $|\mathcal{S}_4^\phi - \mathcal{S}_5^\phi|$ ; i)  $|\mathcal{S}_5^\phi - \mathcal{S}_1^\phi|$ ; and b)  $\epsilon - IS(\mathcal{S}_1^\phi, \mathcal{S}_2^\phi)$ ; d)  $\epsilon - IS(\mathcal{S}_2^\phi, \mathcal{S}_3^\phi)$ ; f)  $\epsilon - IS(\mathcal{S}_3^\phi, \mathcal{S}_4^\phi)$ ; h)  $\epsilon - IS(\mathcal{S}_4^\phi, \mathcal{S}_5^\phi)$ ; j)  $\epsilon - IS(\mathcal{S}_5^\phi, \mathcal{S}_1^\phi)$ .

634 in Figure 18: a)  $|\mathcal{S}_1^\phi - \mathcal{S}_2^\phi|$ ; c)  $|\mathcal{S}_2^\phi - \mathcal{S}_3^\phi|$ ; e)  $|\mathcal{S}_3^\phi - \mathcal{S}_4^\phi|$ ; g)  $|\mathcal{S}_4^\phi - \mathcal{S}_5^\phi|$ ; and i)  $|\mathcal{S}_5^\phi -$   
635  $\mathcal{S}_1^\phi|$ . The second column of Figure 18 shows the  $\epsilon$ -itinerary synchronizations  
636 between the  $i$ -th node and the  $(i-1)$ -th node: b)  $\epsilon - IS(\mathcal{S}_1^\phi, \mathcal{S}_2^\phi)$ ; d)  
637  $\epsilon - IS(\mathcal{S}_2^\phi, \mathcal{S}_3^\phi)$ ; f)  $\epsilon - IS(\mathcal{S}_3^\phi, \mathcal{S}_4^\phi)$ ; h)  $\epsilon - IS(\mathcal{S}_4^\phi, \mathcal{S}_5^\phi)$ ; and j)  $\epsilon - IS(\mathcal{S}_5^\phi, \mathcal{S}_1^\phi)$ .  
638 In conclusion, in this example all the nodes of the ring topology network are  
639 fulfilled the definition of  $\epsilon$ -itinerary synchronization by considering  $\epsilon = 0.02$ ,  
640 see Figure 19. This Figure shows the different curves computed by (16) for  
641  $\epsilon$ -itinerary synchronization between the nodes of a nearly identical network (22)  
642 with coupling strength  $c = 10$  and  $\Gamma = \text{diag}\{1, 0, 0\}$ ; the scroll-degree and initial  
643 condition for each node are given in Table (1): blue line for  $\epsilon - IS(\mathcal{S}_1^\phi, \mathcal{S}_2^\phi)$ ; red  
644 line for  $\epsilon - IS(\mathcal{S}_2^\phi, \mathcal{S}_3^\phi)$ ; black line for  $\epsilon - IS(\mathcal{S}_3^\phi, \mathcal{S}_4^\phi)$ ; green line for  $\epsilon - IS(\mathcal{S}_4^\phi, \mathcal{S}_5^\phi)$ ;  
645 and magenta line for  $\epsilon - IS(\mathcal{S}_5^\phi, \mathcal{S}_1^\phi)$ . However they do not satisfy the definition  
646 of complete synchronization due to the third node are oscillating in a different  
647 manner, see Figure 17 e).

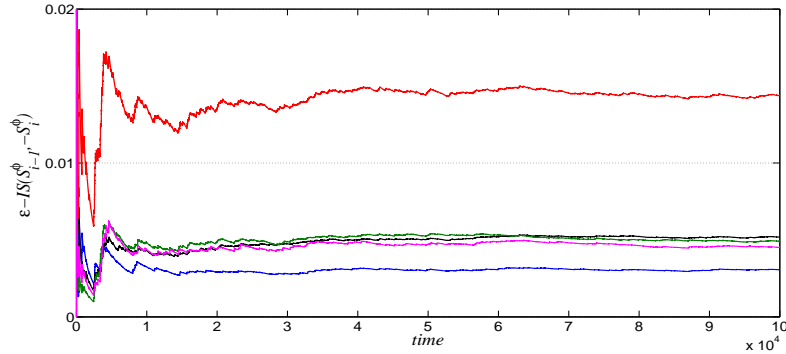


Figure 19: Different curves computed by (16) for  $\epsilon$ -itinerary synchronization between the nodes of a nearly identical network (22) with coupling strength  $c = 10$  and  $\Gamma = \text{diag}\{1, 0, 0\}$ ; the scroll-degree and initial condition for each node are given in Table (1): blue line for  $\epsilon - IS(\mathcal{S}_1^\phi, \mathcal{S}_2^\phi)$ ; red line for  $\epsilon - IS(\mathcal{S}_{s2}^\phi, \mathcal{S}_3^\phi)$ ; black line for  $\epsilon - IS(\mathcal{S}_3^\phi, \mathcal{S}_4^\phi)$ ; green line for  $\epsilon - IS(\mathcal{S}_4^\phi, \mathcal{S}_5^\phi)$ ; and magenta line for  $\epsilon - IS(\mathcal{S}_5^\phi, \mathcal{S}_1^\phi)$ .

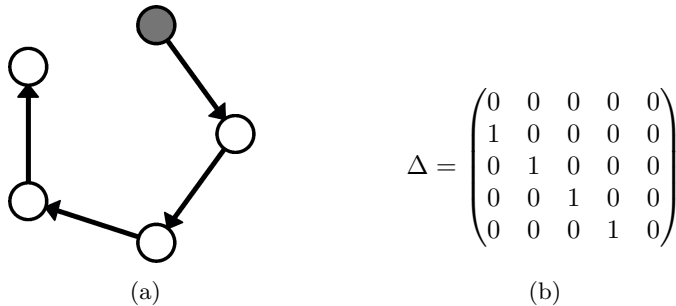


Figure 20: A network of  $N = 5$  nodes coupled in a open ring topology with directional links. (a) The network topology and (b) the coupling matrix.

### 648 7.3 Dynamics in a directed chain topology

649 In this subsection we present numerical results for the case in which the network  
 650 has a directed chain topology. This change transforms the network topology  
 651 from a ring configuration to a chain (open ring) configuration as we illustrate in  
 652 Figure 20; where we also show the corresponding coupling matrix that describes  
 653 this network.

654 After removing a node, the black node in Figure 20 (a), which we call the  
 655 leader node, plays the role of the master system for the rest of the nodes. The  
 656 second node is the slave system for the leader node, but it is also the master  
 657 system for the third node, and so on. The idea is to explore if such a leader  
 658 node governs or not the collective dynamics of the rest of the nodes. In this  
 659 work we assume that the scroll-degree of the master node corresponds to the  
 660 largest or the smallest scroll-degree. Specifically we consider two examples: the



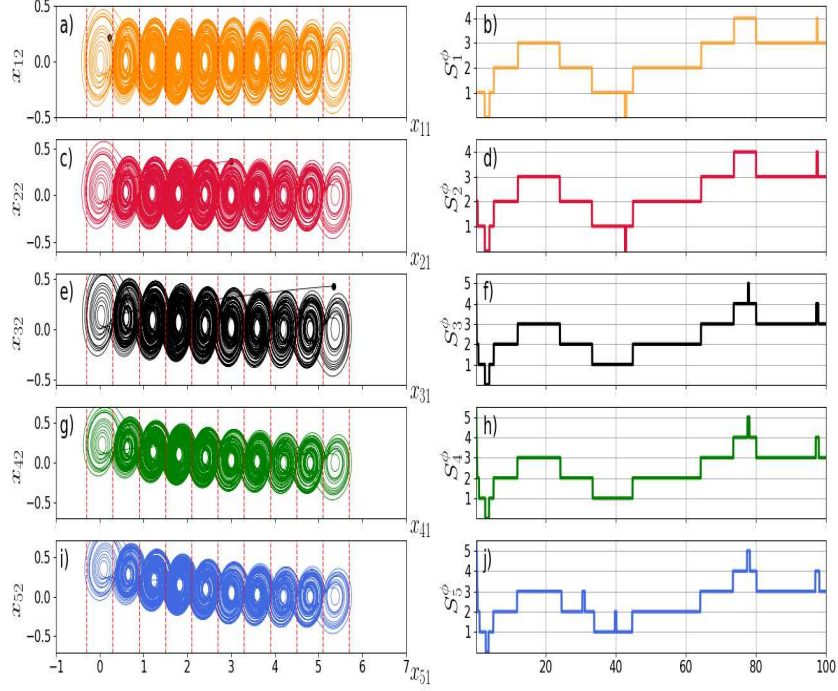


Figure 21: a), c), e), g), i): the projections of the attractors onto the plane  $(x_{i1}, x_{i2})$ , with  $i = 1, \dots, 5$ , of the nodes of a nearly identical network (22) in an directed chain topology with coupling strength  $c = 10$ ,  $\Gamma = \text{diag}\{1, 1, 1\}$  and where the first node has scroll-degree  $\eta_i = 10$ . b), d), f), h), j): the itinerary of each node.

661 first node has scroll-degree ten or three.

### 662 7.3.1 Master system with maximum scroll-degree

663 Figure 21 shows the projections onto the plane  $(x_{i1}, x_{i2})$ , with  $i = 1, \dots, 5$ , of  
 664 the attractors generated in each node by the nearly identical network (22) with a  
 665 chain configuration. For this example we assume that the first node has scroll-  
 666 degree  $\eta_1 = 10$ , and the nodes are connected with coupling strength  $c = 10$   
 667 and inner coupling matrix  $\Gamma = \text{diag}\{1, 1, 1\}$ . The node's scroll-degree and its  
 668 corresponding initial condition are given in Table (1). All the nodes imitate the  
 669 dynamics of the master system and change their dynamics to attain the same  
 670 scroll-degree. In this context, the scroll-degree of the leader node dominates  
 671 and itinerary synchronization is achieved in short periods of time as is shown  
 672 in the second column of Figure 21. However for a long period of time it is  
 673 possible to observe spikes and all the nodes of the network present  $\epsilon$ - Itinerary  
 674 Synchronization for  $\epsilon = 0.02$ , see Figure 22 a). This figure 22 a) shows the

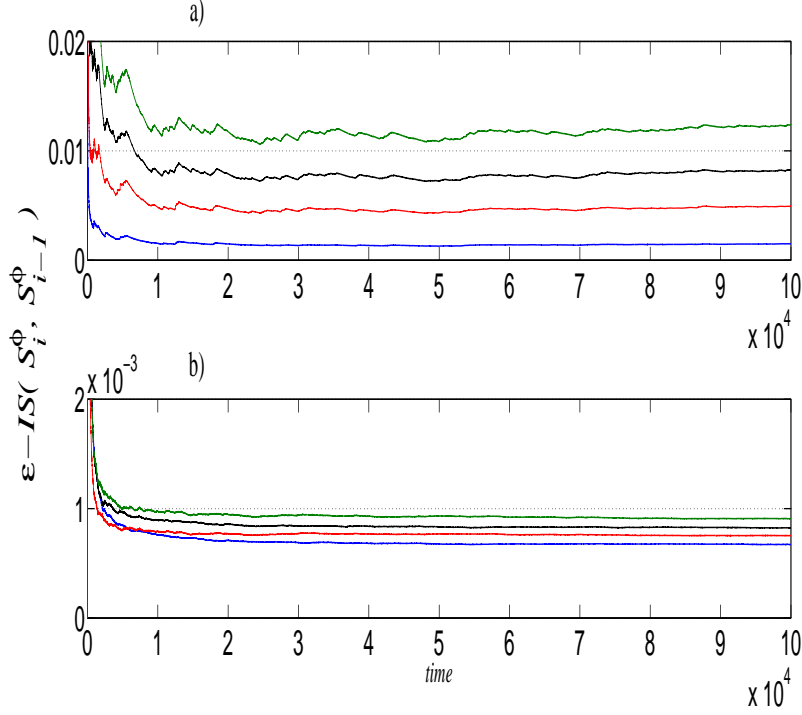


Figure 22: Different curves computed by (16) for  $\epsilon$ -itinerary synchronization between the nodes of a nearly identical network (22) in an directed chain topology with coupling strength  $c = 10$  and  $\Gamma = \text{diag}\{1, 1, 1\}$ ; the scroll-degree and initial condition for each node are given in Table (1): blue line for  $\epsilon - IS(\mathcal{S}_1^\phi, \mathcal{S}_2^\phi)$ ; red line for  $\epsilon - IS(\mathcal{S}_2^\phi, \mathcal{S}_3^\phi)$ ; black line for  $\epsilon - IS(\mathcal{S}_3^\phi, \mathcal{S}_4^\phi)$ ; and green line for  $\epsilon - IS(\mathcal{S}_4^\phi, \mathcal{S}_5^\phi)$ . a) Master system with maximum scroll-degree, and b) Master system with minimum scroll-degree.

675 different curves computed by (16) for  $\epsilon$ -itinerary synchronization between the  
676 nodes of a nearly identical network (22): blue line for  $\epsilon - IS(\mathcal{S}_1^\phi, \mathcal{S}_2^\phi)$ ; red line for  
677  $\epsilon - IS(\mathcal{S}_2^\phi, \mathcal{S}_3^\phi)$ ; black line for  $\epsilon - IS(\mathcal{S}_3^\phi, \mathcal{S}_4^\phi)$ ; and green line for  $\epsilon - IS(\mathcal{S}_4^\phi, \mathcal{S}_5^\phi)$ .

### 678 7.3.2 Master system with minimum scroll-degree

679 Now we assume that after removing the link, the first node has scroll-degree  
680  $\eta_1 = 3$ , and the rest of the nodes have the scroll-degree and initial condition  
681 given in Table (1). As before, we select a coupling strength  $c = 10$  and  $\Gamma =$   
682  $\text{diag}\{1, 1, 1\}$ . In Figure 23 we observe that all the nodes reduce their scroll-  
683 degree to three *i.e.* the nodes adopt the scroll-degree of the first node, in this  
684 case Figure 23 e) shows the leader node. Furthermore, the rest of the nodes  
685 achieve  $\epsilon$ -Itinerary Synchronization for the set of given initial conditions and  
686  $\epsilon = 0.001$ , see Figure 22 b). This figure 22 b) shows the different curves

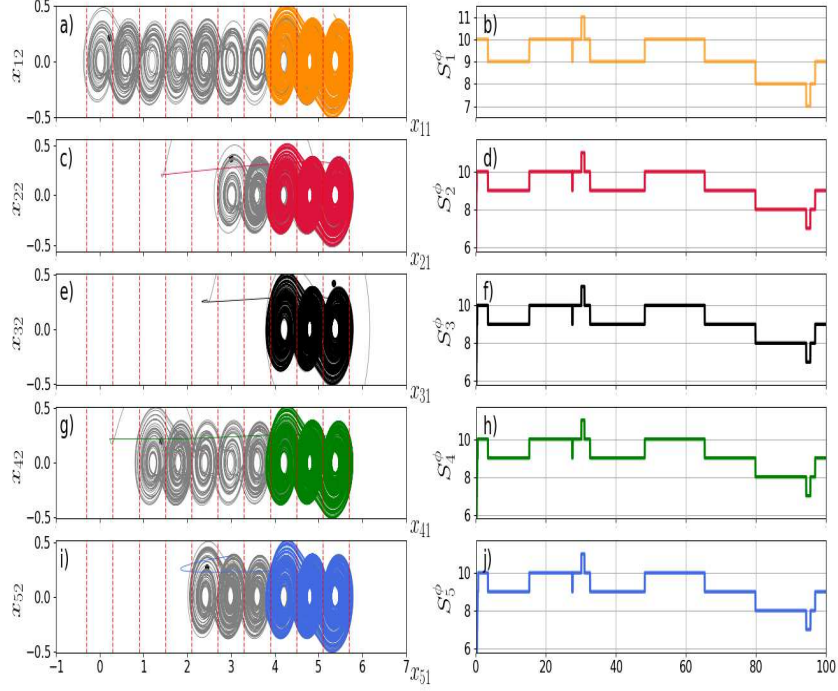


Figure 23: a), c), e), g), i): The projections of the attractors onto the plane  $(x_{i1}, x_{i2})$ , with  $i = 1, \dots, 5$ , of a nearly identical network (22) in a directed chain topology with coupling strength  $c = 10$ ,  $\Gamma = \text{diag}\{1, 1, 1\}$  and where the first node has scroll-degree  $\eta_i = 3$ . b), d), f), h), j): The itinerary of each node.

687 computed by (16) for  $\epsilon$ -itinerary synchronization between the nodes of a nearly  
 688 identical network (22): blue line for  $\epsilon - IS(\mathcal{S}_1^\phi, \mathcal{S}_2^\phi)$ ; red line for  $\epsilon - IS(\mathcal{S}_2^\phi, \mathcal{S}_3^\phi)$ ;  
 689 black line for  $\epsilon - IS(\mathcal{S}_3^\phi, \mathcal{S}_4^\phi)$ ; and green line for  $\epsilon - IS(\mathcal{S}_4^\phi, \mathcal{S}_5^\phi)$ .

## 690 8 Conclusions

691 We have considered PWL systems, generated via heteroclinic orbits and whose  
 692 dynamics exhibits a double scroll attractor. The concept of scroll-degree has  
 693 been introduced to describe the number of scrolls that the PWL system dis-  
 694 plays in its attractor. We study the dynamics of this PWL system by symbolic  
 695 dynamics, given by a natural partition of the state space. Synchronization  
 696 phenomena has been studied in a master-slave system using two inner link-  
 697 ing matrices:  $\Gamma = \text{diag}\{0, 1, 0\}$  and  $\Gamma = \text{diag}\{1, 1, 1\}$ . For both inner linking  
 698 matrices and the coupling strength  $c = 10$  we found that the master-slave sys-  
 699 tem presents itinerary synchronization when the systems are identical, and for

700  $\Gamma = \text{diag}\{0, 1, 0\}$  and the same coupling strength then the master-slave system  
701 presents  $\epsilon$ -itinerary synchronization when the systems are quasi-symmetrical.  
702 This leads to multistability behavior if the scroll-degree of the master system is  
703 less than the slave system.

704 Our numerical results show that for sufficiently large coupling strength,  $\epsilon$ -  
705 itinerary synchronization for small  $\epsilon$  is achieved for different configurations of the  
706 inner coupling matrix. Furthermore, we observe in the multistability regimen  
707 that if the scroll-degree of the master system is less than the slave degree, then  
708 the slave system reduces its scroll-degree and, depending on its initial condition,  
709 it evolves between distinct basins of attractions. On the hand, if the scroll-  
710 degree of the master system is greater than the slave, we observe that the slave  
711 system increase its scroll-degree to be the same as the master, and  $\epsilon$ -itinerary  
712 synchronization is also achieved.

713 The concept of network of nearly identical nodes was introduced to character-  
714 ize a dynamical network composed of PWL systems with different scroll degrees.  
715 We investigated the collective dynamics of an  $N$ -coupled PWL-systems with  
716 different scroll-degree and connected in a master-slave scheme, that is, a unidi-  
717 rectional ring topology. For a network of  $N$ -coupled PWL-systems, we observe  
718 that the node with the smallest scroll-degree governs the collective itinerary of  
719 the network, i.e., the dominant node in a ring configuration network is that  
720 with smallest scroll-degree. Furthermore, we show that the network can display  
721 several behaviors depending on the inner linking matrix  $\Gamma$ . Next, we extend our  
722 results to the case in which we remove a link from the network, transforming  
723 its topology to a directed chain topology. Here we explore two scenarios: the  
724 first node in the chain has the largest scroll-degree, or it has the smallest one.  
725 In the first scenario, we observe that all the nodes increase their scroll-degree  
726 and  $\epsilon$ -itinerary synchronization for small  $\epsilon$  is achieved. For the second scenario  
727 we observe that all the nodes reduce scroll-degree and evolve in the same basin  
728 of attraction of the master system.

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